# The Continuous Classical Boundary Optimal Control of a Couple Nonlinear Parabolic Partial Differential Equations 

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#### Abstract

In this paper the continuous classical boundary optimal control problem of a couple nonlinear partial differential equations of parabolic type is studied. The Galerkin method is used to prove the existence and uniqueness theorem of the state vector solution of a couple nonlinear parabolic partial differential equations for given (fixed) continuous classical boundary control vector. The theorem of the existence of a continuous classical optimal boundary control vector associated with the couple of nonlinear parabolic partial differential equations is proved. The existence of a unique vector solution of the adjoint equations is studied. The Fréchet derivative is derived; Finally The Kuhn-Tucker-Lagrange multipliers theorems is developed and then is used to prove the necessary conditions theorem and the sufficient conditions theorem of optimality of a couple of nonlinear parabolic equations with equality and inequality constraints. [DOI: 10.22401/ANJS.00.1.17]


Keywords: boundary optimal control, couple nonlinear parabolic partial differential equations.

## 1. Introduction

The objective of optimal control theory is to determine the control signals that will cause a process to satisfy the physical constraints and at the same time minimize (or maximize) some performance criterion, [8]. Control theory is an application-oriented mathematics that deals with the basic principles underlying the analysis and design of (control) system. Systems can be engineering (air conditioner, air craft, and CD player etc.), economic, and biological, [12]. In general, there are many optimal control problems are governed either by ODEs as Orpel in 2009[11] or by different types of PDEs and are subject to control and state constraints, as El-Borari and et al in 2013 [9], and Wang, Y. and et al in 2015 [15], which are studied an optimal control of parabolic partial differential equations, Farag, M. H. in 2014[10] studied classical optimal control of hyperbolic partial differential equations, Diaz and et al in 2012 [7] studied a optimal control of elliptic partial differential equations, Al-Rawdanee, E. in 2014 [3] studied an a classical optimal control of a coupled of nonlinear elliptic partial differential equations and M. K. Ghufran in 2016 [4] studied a classical optimal control of a coupled of nonlinear parabolic partial differential equations while, Al-Hawasy, J. in 2016 [2] studied a classical optimal control of a coupled
of nonlinear hyperbolic partial differential equations.

This paper deals with, the proof of the existence and uniqueness theorem of the state vector solution of a couple nonlinear parabolic partial differential equations where the continuous classical boundary control vector is given, the existence theorem of a continuous classical boundary optimal control vector associated with a couple nonlinear partial differential equations of parabolic type is proved, also the derivation of the Fréchet derivative is done, the study of the existence and uniqueness of the vector solution of the adjoint equations which corresponds to the state vector. Finally, the Kuhn-TuckerLagrange multipliers theorem is developed and is used to prove the necessary conditions theorem and the sufficient conditions theorem of optimality of a couple of nonlinear parabolic equations with equality and inequality constraints.

## 2. Description of the Problem

Let $I=(0, T), T<\infty, \Omega \subset \mathbb{R}^{2}$ be an open and bounded region with Lipschitz boundary $\Gamma=\partial \Omega, Q=\Omega \times I, \Sigma=\Gamma \times I$. Consider the following continuous boundary optimal control problem:

The state equation is given by the following nonlinear parabolic equation:
$y_{1 t}-\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(a_{i j}(x, t) \frac{\partial y_{1}}{\partial x_{j}}\right)+b_{1}(x, t) y_{1}-$ $b(x, t) y_{2}=f_{1}\left(x, t, y_{1}\right)$, in Q
$y_{2 t}-\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(b_{i j}(x, t) \frac{\partial y_{2}}{\partial x_{j}}\right)+b_{2}(x, t) y_{2}+$
$b(x, t) y_{1}=f_{2}\left(x, t, y_{2}\right)$, in Q
$\sum_{i, j=1}^{n} a_{i j} \frac{\partial y_{1}}{\partial n}=u_{1}(x, t)$, on $\Sigma$
$y_{1}(x, 0)=y_{1}^{0}(x)$, on $\Omega$
$\sum_{i, j=1}^{n} b_{i j} \frac{\partial y_{2}}{\partial n}=u_{2}(x, t)$, on $\Sigma$
$y_{2}(x, 0)=y_{2}^{0}(x)$, on $\Omega$
where for all $x=\left(x_{1}, x_{2}\right), \quad\left(y_{1}, y_{2}\right) \in$ $\left(H^{1}(Q)\right)^{2}$ is the state vector, $\left(u_{1}, u_{2}\right) \in$ $\left(L^{2}(\Sigma)\right)^{2}$ is the classical boundary control vector, $\left(f_{1}, f_{2}\right) \in\left(L^{2}(Q)\right)^{2}$ is a vector of a given function defined $(Q \times \mathbb{R}) \times(Q \times \mathbb{R})$, and $a_{i j}(x, t), b_{i j}(x, t), b(x, t)$ and $b_{i}(x, t) \in$ $C^{\infty}(Q)$.
$\vec{W}_{A}=$
$\left\{\vec{w} \in L^{2}(\Sigma) \times L^{2}(\Sigma) \mid \vec{w} \in \vec{U}\right.$ a.e. in $\left.\Sigma, G_{1}(\vec{w})=0, G_{2}(\vec{w}) \leq 0\right\}$ $\vec{U} \subset \mathbb{R}^{2}$.
The cost function is
$G_{0}(\vec{u})=$
$\int_{Q}\left[g_{01}\left(x, t, y_{1}\right)+\right.$
$\left.g_{02}\left(x, t, y_{2}\right)\right] d x d t+\int_{\Sigma}\left[h_{01}\left(x, t, u_{1}\right)+h_{02}\left(x, t, u_{2}\right)\right] d \sigma$
The constraints on the state and the control vectors are
$G_{1}(\vec{u})=\int_{Q}\left[g_{11}\left(x, t, y_{1}\right)+g_{12}\left(x, t, y_{2}\right)\right] d x d t+$
$\int_{\Sigma}\left[h_{11}\left(x, t, u_{1}\right)+h_{12}\left(x, t, u_{2}\right)\right] d \sigma=0$
$G_{2}(\vec{u})=\int_{Q}\left[g_{21}\left(x, t, y_{1}\right)+g_{22}\left(x, t, y_{2}\right)\right] d x d t+$
$\int_{\Sigma}\left[h_{21}\left(x, t, u_{1}\right)+h_{22}\left(x, t, u_{2}\right)\right] d \sigma \leq 0$
where $\left(y_{1}, y_{2}\right)=\left(y_{u_{1}}, y_{u_{2}}\right)$ is the solution of (1-6) corresponding to the boundary control vector $\left(u_{1}, u_{2}\right)$.
$\operatorname{Let} \vec{V}=V \times V=\left\{\vec{v}: \vec{v} \in\left(H^{1}(\Omega)\right)^{2}\right\}, \vec{v}=\left(v_{1}, v_{2}\right)$.
We denote by $(v, v)_{\Omega}$ and $\|v\|_{0}$ (by $(v, v)_{\Gamma}$ and $\|v\|_{\Gamma}$ ) the inner product and the norm in $\mathrm{L}^{2}(\Omega)\left(\right.$ in $\mathrm{L}^{2}(\Gamma)$ ), by $(v, v)_{1}$ and $\|v\|_{1}$ the inner product and the norm in $H^{1}(\Omega)$, by $(\vec{v}, \vec{v})_{\Omega}$ and $\|\vec{v}\|_{0}$ (by $(\vec{v}, \vec{v})_{\Gamma}$ and $\|\vec{v}\|_{\Gamma}$ ) the inner product and the norm in $L^{2}(\Omega) \times L^{2}(\Omega)$ ( in $\left.L^{2}(\Gamma) \times L^{2}(\Gamma)\right) \quad$ by $\quad(\vec{v}, \vec{v})_{1}=\left(v_{1}, v_{1}\right)_{1}+$ $\left(v_{2}, v_{2}\right)_{1}$ and $\|\vec{v}\|_{1}^{2}=\left\|v_{1}\right\|_{1}^{2}+\left\|v_{2}\right\|_{1}^{2} \quad$ the inner product and the norm in $\vec{V}$ and $\vec{V}^{*}$ is the dual of $\vec{V}$.

## 3. Weak Formulation of the State Equations

The weak forms of the problem (1-6) when $\vec{y} \in\left(H_{0}^{1}(\mathrm{Q})\right)^{2}$ are given almost everywhere on $I\left(\forall v_{1}, v_{2} \in V, y_{1}(., t), y_{2}(., t) \in V\right)$ by $\left\langle y_{1 t}, v_{1}\right\rangle+a_{1}\left(t, y_{1}, v_{1}\right)+\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-$ $\left(b(t) y_{2}, v_{1}\right)_{\Omega}=\left(f_{1}, v_{1}\right)_{\Omega}+\left(u_{1}, v_{1}\right)_{\Gamma}$,
$\left(y_{1}^{0}, v_{1}\right)_{\Omega}=\left(y_{1}(0), v_{1}\right)_{\Omega}$
and
$\left\langle y_{2 t}, v_{2}\right\rangle+a_{2}\left(t, y_{2}, v_{2}\right)+\left(b_{2}(t) y_{2}, v_{2}\right)_{\Omega}+$
$\left(b(t) y_{1}, v_{2}\right)_{\Omega}=\left(f_{2}, v_{2}\right)_{\Omega}+\left(u_{2}, v_{2}\right)_{\Gamma}$,
$\left(y_{2}^{0}, v_{2}\right)_{\Omega}=\left(y_{2}(0), v_{2}\right)_{\Omega}$
Where $a_{1}\left(t, y_{1}, v_{1}\right)=\int_{\Omega} \sum_{i, j=1}^{n} a_{i j} \frac{\partial y_{1}}{\partial x_{i}} \frac{\partial v_{1}}{\partial x_{j}} d x$, $a_{2}\left(t, y_{2}, v_{2}\right)=\int_{\Omega} \sum_{i, j=1}^{n} b_{i j} \frac{\partial y_{2}}{\partial x_{i}} \frac{\partial v_{2}}{\partial x_{j}} d x$.
To study the existence of unique solution of the weak form (10-11), we consider the following assumption.

## Assumptions (A):

(i) $f_{i}$ is of a Carathéodory type on $Q \times \mathbb{R}$, satisfies the following sub linearity condition for $y_{i}$, i.e. $\left|f_{i}\left(x, t, y_{i}\right)\right| \leq$ $\eta_{i}(x, t)+c_{i}\left|y_{i}\right|$

Where $(x, t) \in Q, y_{i} \in \mathbb{R}, c_{i}>0$ and $\eta_{i} \in L^{2}(Q, \mathbb{R}), \forall i=1,2$
(ii) $f_{i}$ is Lipschitz w.r.t. $y_{i}$, i.e. $\mid f_{i}\left(x, t, y_{i}\right)-$ $f_{i}\left(x, t, \hat{y}_{i}\right)\left|\leq L_{i}\right| y_{i}-\hat{y}_{i} \mid$
(iii) Where $(x, t) \in Q, y_{i}, \hat{y}_{i} \in \mathbb{R}$ and $L_{i}>0$, $\forall i=1,2$
(iv) $c(t, \vec{y}, \vec{y})=$ $a_{1}\left(t, y_{1}, y_{1}\right)+\left(b_{1}(t) y_{1}, y_{1}\right)_{\Omega}+$
$a_{2}\left(t, y_{2}, y_{2}\right)+\left(b_{2}(t) y_{2}, y_{2}\right)_{\Omega}, \quad$ and $|c(t, \vec{y}, \vec{y})| \leq \alpha\|\vec{y}\|_{1}\|\vec{v}\|_{1}, c(t, \vec{y}, \vec{y}) \geq$ $\bar{\alpha}\|\vec{y}\|_{1}^{2}$, where $\alpha, \bar{\alpha}$ are real positive constants

## Proposition (3.1), [6]:

Let $f: \Omega \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{m}$ is of a Carathéodory type, let $F$ be a functional, such that
$F(y)=\int_{\Omega} f(x, y(x)) d x$,
where $\Omega$ is measurable subset of $\mathbb{R}^{d}(d=2,3)$, and suppose that
$\|f(x, y)\| \leq \zeta(x)+\eta(x)\|y\|^{\alpha}, \forall(x, y) \in \Omega \times$ $\mathbb{R}^{n}, y \in L^{p}\left(\Omega \times \mathbb{R}^{n}\right)$
Where $\zeta(x) \in L^{1}(\Omega \times \mathbb{R}), \eta \in L^{\frac{p}{p-\alpha}}(\Omega \times \mathbb{R})$, and $\alpha \in[0, p]$, if $p \in[1, \infty)$, and $\eta \equiv 0$, if $p=\infty$. Then $F$ is continuous on $L^{p}\left(\Omega \times \mathbb{R}^{n}\right)$.

Proposition (3.2), [6]:
Let $f \& f_{y}: \Omega \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{m}$ are Carathéodory type, let $F: L^{p}(\Omega) \rightarrow \mathbb{R}$ be a functional, such that, $F(y)=\int_{\Omega} f_{y}(x, y(x)) d x$ where $\Omega$ is a measurable subset of $\mathbb{R}^{d}$, and
$\left\|f_{y}(x, y)\right\| \leq \zeta(x)+\eta(x)\|y\|^{\frac{\beta}{q}}, \quad \forall(x, y) \in$ $\Omega \times \mathbb{R}^{n}$,
Where $\quad \zeta \in L^{q}(\Omega \times \mathbb{R}), \quad \frac{1}{p}+\frac{1}{q}=1, \quad \eta \in$ $L^{\frac{p q}{p-\beta}}(\Omega \times \mathbb{R}), \beta \in[0, p]$ if $p \neq \infty$, and $\eta \equiv 0$, if $p=\infty$

Then the Fréchet derivative of $F$ exists for each $y \in L^{p}\left(\Omega \times \mathbb{R}^{n}\right)$ and is given by:
$\mathrm{G}(y) h=\int_{\Omega} f_{y}(x, y(x)) h(x) d x$.

## Theorem (3.1), [14]:

Let $D$ be a measurable subset of $\mathbb{R}^{d}$, $\emptyset: D \rightarrow \mathbb{R}$ and $\emptyset \in L^{1}(D, \mathbb{R})$.

If the following inequality is satisfied $\int_{S} \phi(v) d v \geq 0 \quad$ (or $\quad \leq 0,=0$ ), for each measurable set $S \subset D$, then $\phi(v) \geq 0($ or $\leq 0,=0)$, a.e. in $D$.

Theorem (3.2) (Existence and Uniqueness of
Solution of the State Equations):
With assumptions (A), for each fixed boundary control $\vec{u} \in\left(L_{2}(\Sigma)\right)^{2}$, the weak form of the state equations (10-11) has a unique solution $\vec{y}=\left(y_{1}, y_{2}\right)$,
s.t.
$\vec{y} \in\left(L^{2}(I, V)\right)^{2}$ and $\vec{y}_{t}=\left(y_{1 t}, y_{2 t}\right) \in$
$\left(L^{2}\left(I, V^{*}\right)\right)^{2}$
Proof:
Let $\vec{V}_{n} \subset \vec{V}$ be the set of continuous and piecewise affine functions in $\Omega$, let $\left\{\vec{v}_{1}, \vec{v}_{2}, \ldots, \vec{v}_{n}\right\}$ be basis of $\vec{V}_{n}$ where $n=2 N$ (where $N$ is the dimension of each $V$ ), then the approximate solution $\vec{y}$ of $(10-11)$ is approximated by $\vec{y}_{n}=\left(y_{1 n}, y_{2 n}\right)$, such that, for each $n$
$\vec{y}_{n}=\sum_{j=1}^{n} C_{j}(t) \vec{v}_{j}(x)$
where
$\vec{v}_{j}=\left((2-\ell) v_{1 k},(\ell-1) v_{2 k}\right)$, and $C_{j}=c_{l j}$
For $j=k+n(\ell-1), \quad k=1, \ldots N, \ell=1,2$, and $c_{l j}(t)$ is unknown function of t .

The weak forms of the state equations (10) and (11) can be approximated w.r.t. the space variable, using the Galerkin's method to get,
$\left\langle y_{1 n t}, v_{1}\right\rangle+a_{1}\left(t, y_{1 n}, v_{1}\right)+$
$\left(b_{1}(t) y_{1 n}, v_{1}\right)_{\Omega}-\left(b(t) y_{2 n}, v_{1}\right)_{\Omega}=$
$\left(f_{1}\left(y_{1 n}\right), v_{1}\right)_{\Omega}+\left(u_{1}, v_{1}\right)_{\Gamma}$
$\left(y_{1 n}^{0}, v_{1}\right)_{\Omega}=\left(y_{1}^{0}, v_{1}\right)_{\Omega}, \forall v_{1} \in V_{n}$
and
$\left\langle y_{2 n t}, v_{2}\right\rangle+a_{2}\left(t, y_{2 n}, v_{2}\right)+$
$\left(b_{2}(t) y_{2 n}, v_{2}\right)_{\Omega}+\left(b(t) y_{1 n}, v_{2}\right)_{\Omega}=$
$\left(f_{2}\left(y_{2 n}\right), v_{2}\right)_{\Omega}+\left(u_{2}, v_{2}\right)_{\Gamma}$
$\left(y_{2 n}^{0}, v_{2}\right)_{\Omega}=\left(y_{2}^{0}, v_{2}\right)_{\Omega}$,
where $y_{\text {in }}^{0}=y_{\text {in }}(x, 0) \in V_{n} \subset V \subset L^{2}(\Omega)$ is
the projection of $y_{i}^{0}$ for the norm $\|$. $\|_{0, \text { i.e. }}$
$\left(y_{i n}^{0}, v_{i}\right)_{\Omega}=\left(y_{i}^{0}, v_{i}\right)_{\Omega} \forall v_{i} \in V_{n} \Leftrightarrow \| y_{i n}^{0}-$
$y_{i}^{0}\left\|_{0} \leq\right\| y_{i}^{0}-v_{i} \|_{0}, \forall v_{i} \in V_{n}, \forall i=1,2$
By substituting (12) in (13 a\& b) and in (14 a \& b), one obtains
$A_{1} C_{1}^{\prime}(t)+D_{1} C_{1}(t)-E_{1} C_{2}(t)=$
$b_{1}\left(\bar{V}_{1}^{T}(x) C_{1}(t)\right)$
$A_{1} C_{1}(0)=b_{1}^{0}$
$A_{2} C_{2}^{\prime}(t)+D_{2} C_{2}(t)+E_{2} C_{1}(t)=$
$b_{2}\left(\bar{V}_{2}^{T}(x) C_{2}(t)\right)$.
$B C_{2}(0)=b_{2}^{0}$
where:
$A_{1}=\left(a_{i j}\right)_{n \times n}, a_{i j}=\left(v_{1 j}, v_{1 i}\right)_{\Omega}$,
$D_{1}=\left(d_{i j}\right)_{n \times n}, d_{i j}=\left[a_{1}\left(t, v_{1 j}, v_{1 i}\right)+\right.$
$\left.\left(b_{1}(t) v_{1 j}, v_{1 i}\right)_{\Omega}\right], E_{1}=\left(e_{i j}\right)_{n \times n}, e_{i j}=$
$\left(b(t) v_{2 j}, v_{1 i}\right)_{\Omega}, C_{\ell}(t)=\left(c_{\ell j}(t)\right)_{n \times 1}$,
$C_{\ell}^{\prime}(t)=\left(c_{\ell j}^{\prime}(t)\right)_{n \times 1}, C_{\ell}(0)=\left(c_{\ell j}(0)\right)_{n \times 1}$,
$b_{\ell}=\left(b_{l i}\right)_{n \times 1}, b_{\ell i}=\left(f_{\ell}\left(\bar{v}_{\ell}^{T} C_{\ell}(t)\right), v_{\ell i}\right)_{\Omega}+$
$\left(u_{\ell}, v_{\ell i}\right)_{\Gamma}, \bar{v}_{\ell}=\left(v_{\ell}\right)_{n \times 1}, b_{\ell}^{0}=\left(b_{\ell i}^{0}\right), b_{\ell i}^{0}=$
$\left(y_{\ell}^{0}, v_{\ell i}\right)_{\Omega}$, and $A_{2}=\left(b_{i j}\right)_{n \times n}, b_{i j}=$
$\left(v_{2 j}, v_{2 i}\right)_{\Omega}, D_{2}=\left(f_{i j}\right)_{n \times n}, f_{i j}=$
$\left[a_{2}\left(t, v_{2 j}, v_{2 i}\right)+\left(b_{2}(t) v_{2 j}, v_{2 i}\right)_{\Omega}\right], E_{2}=$
$\left(h_{i j}\right)_{n \times n}, h_{i j}=\left(b(t) v_{1 i}, v_{2 i}\right)_{\Omega}, \ell=1,2$.
From assumptions (A), easily once can get that the matrices $A_{1} \& A_{2}$ are positive definite, therefore the system $\left(12^{\prime}-13^{\prime}\right)$ of $1^{\text {st }}$ order differential equation has unique solution [5].

Now, to show the norm $\left\|\vec{y}_{n}^{0}\right\|_{0}$ is bounded: Since $\vec{y}^{0} \in\left(L^{2}(\Omega)\right)^{2}$, then there exists $\left\{\vec{v}_{n}^{0}\right\}$, with $\vec{v}_{n}^{0} \in \vec{V}_{n}$ such that $\vec{v}_{n}^{0} \rightarrow \vec{y}^{0}$ strongly in $\left(L^{2}(\Omega)\right)^{2}$ and then from the projection theorem and (13b \& 14b) one obtains that $\vec{y}_{n}^{0} \rightarrow \vec{y}^{0}$
Strongly in $\left(L^{2}(\Omega)\right)^{2}$ with $\left\|\vec{y}_{n}^{0}\right\|_{0} \leq b_{1}$.
The norm $\left\|\vec{y}_{n}(t)\right\|_{L^{\infty}\left(1, L^{2}(\Omega)\right)}$ and $\left\|\vec{y}_{n}(t)\right\|_{Q}$ are bounded:

Setting $v_{1}=y_{1 n}$ in (13a) and $v_{2}=y_{2 n}$ in (14a), integrating both sides of each obtaining equation for $t$ from 0 to $T$, and adding them finally with Assumption (A-iii), one has
$\int_{0}^{T}\left\langle\vec{y}_{n t}, \vec{y}_{n}\right\rangle d t+\int_{0}^{T} c\left(t, \vec{y}_{n}, \vec{y}_{n}\right) d t=$
$\int_{0}^{T}\left(f_{1}\left(y_{1 n}\right), y_{1 n}\right)_{\Omega} d t+\int_{0}^{T}\left(f_{2}\left(y_{2 n}\right), y_{2 n}\right)_{\Omega} d t+$
$\int_{0}^{T}\left(u_{1}, y_{1 n}\right)_{\Gamma} d t+\int_{0}^{T}\left(u_{2}, y_{2 n}\right)_{\Gamma} d t$,
Since $\quad \vec{y}_{n t} \in\left(L^{2}\left(I, V^{*}\right)\right)^{2}=\left(L^{2}(I, V)\right)^{2}$ and $\vec{y}_{n} \in\left(L^{2}(I, V)\right)^{2}$ in the $1^{\text {st }}$ term of the L.H.S. of (15), hence for this term we can use Lemma 1.2 in [13] and since the $2^{\text {nd }}$ term is positive, taking $T=t \in[0, T]$, finally using, and Assumption (A-i) for the $1^{\text {st }}$ two terms in the right hand side (for briefly will use R.H.S. from here and next) of (15), one has
$\int_{0}^{t} \frac{d}{d t}\left\|\vec{y}_{n}(t)\right\|_{0}^{2} d t \leq\left\|\eta_{1}\right\|_{Q}^{2}+\left\|\eta_{2}\right\|_{Q}^{2}+$
$\left\|u_{1}\right\|_{\Sigma}^{2}+\left\|u_{2}\right\|_{\Sigma}^{2}+c_{5} \int_{0}^{t}\left\|\vec{y}_{n}\right\|_{0}^{2} d t$,
$\Rightarrow\left\|\vec{y}_{n}(t)\right\|_{0}^{2}-\left\|\vec{y}_{n}(0)\right\|_{0}^{2} \leq m_{1}+m_{2}+c_{1}+$
$c_{2}+c_{5} \int_{0}^{t}\left\|\vec{y}_{n}\right\|_{0}^{2} d t$
$\Rightarrow\left\|\vec{y}_{n}(t)\right\|_{0}^{2} \leq m^{*}+c_{5} \int_{0}^{t}\left\|\vec{y}_{n}\right\|_{0}^{2} d t$
By using the Belman- Gronwall inequality, one gets
$\Rightarrow\left\|\vec{y}_{n}(t)\right\|_{L^{\infty}\left(I, L^{2}(\Omega)\right)} \leq h_{9}$
hence

$$
\begin{aligned}
\left\|\vec{y}_{n}(t)\right\|_{Q}^{2} & =\int_{0}^{T}\left\|\vec{y}_{n}\right\|_{0}^{2} d t \\
& \leq T \max _{t \in[0, T]}\left\|\vec{y}_{n}(t)\right\|_{0}^{2} \\
& \leq T h_{8}=h_{10}^{2}=h_{10}
\end{aligned}
$$

The norm $\left\|\vec{y}_{n}(t)\right\|_{L^{2}(I, V)}$ is bounded:
Again by using Lemma 1.2 in [13] for the $1^{\text {st }}$ term in the L.H.S. of (15), then using same results which are obtained from the R.H.S., finally setting $t=T$, and $\left\|\vec{y}_{n}(T)\right\|_{0}^{2} \geq 0$, equation (15) becomes
$\left\|\vec{y}_{n}(T)\right\|_{0}^{2}+2 \bar{\alpha} \int_{0}^{T}\left\|\vec{y}_{n}\right\|_{1}^{2} d t \leq\left\|\eta_{1}\right\|_{Q}^{2}+$
$\left\|\eta_{2}\right\|_{Q}^{2}+\left\|u_{1}\right\|_{\Sigma}^{2}+\left\|u_{2}\right\|_{\Sigma}^{2}+c_{5}\left\|\vec{y}_{n}\right\|_{Q}^{2}+$
$\left\|\vec{y}_{n}(0)\right\|_{0}^{2}$
$\Rightarrow\left\|\vec{y}_{n}\right\|_{L^{2}(I, V)} \leq h_{11}$

## The convergence of the solution:

Let $\left\{\vec{V}_{n}\right\}_{n=1}^{\infty}$ be a sequence of subspaces of $\vec{V}$, such that $\forall \vec{v}=\left(v_{1}, v_{2}\right) \in \vec{V}$, there exists a sequence $\left\{\vec{v}_{n}\right\}$ with $\vec{v}_{n}=\left(v_{1 n}, v_{2 n}\right) \in \vec{V}_{n}, \forall n$, and $\quad \vec{v}_{n} \rightarrow \vec{v} \quad$ strongly $\quad$ in $\vec{V} \Rightarrow \vec{v}_{n} \rightarrow \vec{v}$ strongly in $\left(L^{2}(\Omega)\right)^{2}$.

Since for each, with $\vec{V}_{n} \subset \vec{V}$, problems ( $13 \mathrm{a} \& \mathrm{~b}$ ) and ( $14 \mathrm{a} \& \mathrm{~b}$ ) have a unique solutions $y_{1 n}, y_{2 n}$ respectively, hence corresponding to the sequence subspaces $\left\{\vec{V}_{n}\right\}_{n=1}^{\infty}$, one obtain a sequence of approximation problems (13 $\mathrm{a} \& \mathrm{~b}$ ) and ( $14 \mathrm{a} \& \mathrm{~b}$ ), by substituting $\vec{v}=\vec{v}_{n}=$ $\left(v_{1 n}, v_{2 n}\right)$ for $n=1,2, \ldots, \quad$ in these approximation problem, one gets
$\left\langle y_{1 n t}, v_{1 n}\right\rangle+a_{1}\left(t, y_{1 n}, v_{1 n}\right)+$
$\left(b_{1}(t) y_{1 n}, v_{1 n}\right)_{\Omega}-\left(b(t) y_{2 n}, v_{1 n}\right)_{\Omega}=$
$\left(f_{1}\left(y_{1 n}\right), v_{1 n}\right)_{\Omega}+\left(u_{1}, v_{1 n}\right)_{\Gamma}, \quad \forall y_{1 n}, y_{2 n} \in$ $L^{2}\left(I, V_{n}\right)$ a.e in I
$\left(y_{1 n}^{0}, v_{1 n}\right)_{\Omega}=\left(y_{1}^{0}, v_{1 n}\right)_{\Omega}, \forall v_{1 n} \in V_{n}, \forall n$
and
$\left\langle y_{2 n t}, v_{2 n}\right\rangle+$
$a_{2}\left(t, y_{2 n}, v_{2 n}\right)+\left(b_{2}(t) y_{2 n}, v_{2 n}\right)_{\Omega}+$
$\left(b(t) y_{1 n}, v_{2 n}\right)_{\Omega}=$
$\left(f_{2}\left(y_{2 n}\right), v_{2 n}\right)_{\Omega}+\left(u_{2}, v_{2 n}\right)_{\Gamma}, \quad \forall y_{1 n}, y_{2 n} \in$ $L^{2}\left(I, V_{n}\right)$ a.e.in I
$\left(y_{2 n}^{0}, v_{2 n}\right)_{\Omega}=\left(y_{2}^{0}, v_{2 n}\right)_{\Omega}, \forall v_{2 n} \in V_{n}, \forall n$
which has a sequence of solutions $\left\{\vec{y}_{n}\right\}_{n=1}^{\infty}$, where $\mathrm{nm} \vec{y}_{n}=\left(y_{1 n}, y_{2 n}\right)$. Since the norms $\left\|\vec{y}_{n}\right\|_{L^{2}(\boldsymbol{Q})}$ and $\left\|\vec{y}_{n}\right\|_{L^{2}(I, V)}$ are bounded, then by Alaoglu's theorem, there exists a subsequence of $\left\{\vec{y}_{n}\right\}_{n \in N}$, say again $\left\{\vec{y}_{n}\right\}_{n \in N}$ such that $\vec{y}_{n} \rightarrow \vec{y}$ weakly in $\left(L^{2}(Q)\right)^{2}$ and in $\left(L^{2}(I, V)\right)^{2}$.

Then through the First Compactness Theorem, Assumption (A-i), and the bounded norms results, once get $\vec{y}_{n} \rightarrow \vec{y}$ strongly in $\left(L^{2}(Q)\right)^{2}$.

Now, consider the weak state equations ( $16 \mathrm{a} \& \mathrm{~b}$ ), ( $17 \mathrm{a} \& \mathrm{~b}$ ) and take any arbitrary, $v_{1}, v_{2} \in V$, then there exists a sequence $\left\{v_{1 n}\right\}$, $\left\{v_{2 n}\right\}$ respectively, $v_{i n} \in V_{n}, \forall n$, such that $v_{\text {in }} \rightarrow v_{i}$ strongly in $V$ (which gives $v_{\text {in }} \rightarrow v_{i}$ strongly in $\left.L^{2}(\Omega)\right), \forall i=1,2$.

Multiplying both sides of (16a) and (17a) by $\varphi_{i}(t) \in C^{1}[0, T] \quad$ respectively, with $\varphi_{i}(T)=0, \forall i=1,2$, integrating with respect to $t$ from 0 to $T$, and then integrating by parts the $1^{\text {st }}$ term in the L.H.S. of each obtained equation, one gets that
$-\int_{0}^{T}\left(y_{1 n}, v_{1 n}\right) \varphi_{1}^{\prime}(t) d t+$
$\int_{0}^{T}\left[a_{1}\left(t, y_{1 n}, v_{1 n}\right)+\left(b_{1}(t) y_{1 n}, v_{1 n}\right)_{\Omega}-\right.$
$\left.\left(b(t) y_{2 n}, v_{1 n}\right)_{\Omega}\right] \varphi_{1}(t) d t=$
$\int_{0}^{T}\left(f_{1}\left(y_{1 n}\right), v_{1 n}\right)_{\Omega} \varphi_{1}(t) d t+$
$\int_{0}^{T}\left(u_{1}, v_{1}\right)_{\Omega} \varphi_{1}(t) d t+\left(y_{1 n}^{0}, v_{1 n}\right)_{\Omega} \varphi_{1}(0)$
and
$-\int_{0}^{T}\left(y_{2 n}, v_{2 n}\right) \varphi_{2}^{\prime}(t) d t+$
$\int_{0}^{T}\left[a_{2}\left(t, y_{2 n}, v_{2 n}\right)+\left(b_{2}(t) y_{2 n}, v_{2 n}\right)_{\Omega}+\right.$
$\left.\left(b(t) y_{1 n}, v_{2 n}\right)_{\Omega}\right] \varphi_{2}(t) d t=$
$\int_{0}^{T}\left(f_{2}\left(y_{2 n}\right), v_{2 n}\right)_{\Omega} \varphi_{2}(t) d t+$
$\int_{0}^{T}\left(u_{2}, v_{2}\right)_{\Gamma} \varphi_{2}(t) d t+\left(y_{2 n}^{0}, v_{2 n}\right)_{\Omega} \varphi_{2}(0)$
Since $\forall i=1,2, \quad y_{i n} \rightarrow y_{i}$ weakly in $L^{2}(Q), y_{i n}^{0} \rightarrow y_{i}^{0}$ strongly in $L^{2}(\Omega)$, and
$\left.\begin{array}{c}v_{\text {in }} \rightarrow v_{i} \text { strongly in } L^{2}(\Omega) \\ v_{\text {in }} \rightarrow v_{i} \text { strongly in } V\end{array}\right\} \Rightarrow$
$\left\{\begin{array}{c}v_{\text {in }} \varphi_{i}^{\prime} \rightarrow v_{i} \varphi_{i}^{\prime} \text { strongly in } L^{2}(Q) \\ v_{\text {in }} \varphi_{i} \rightarrow v_{i} \varphi_{i} \text { strongly in } L^{2}(I, V)\end{array}\right.$
Then the following convergences hold
$\int_{0}^{T}\left(y_{1 n}, v_{1 n}\right) \varphi_{1}^{\prime}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1 n}, v_{1 n}\right)+\right.$ $\left(b_{1}(t) y_{1 n}, v_{1 n}\right)_{\Omega}-$
$\left.\left(b(t) y_{2 n}, v_{1 n}\right)_{\Omega}\right] \varphi_{1}(t) d t \rightarrow$
$\int_{0}^{T}\left(y_{1}, v_{1}\right) \varphi_{1}^{\prime}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1}, v_{1}\right)+\right.$
$\left.\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-\left(b(t) y_{2}, v_{1}\right)_{\Omega}\right] \varphi_{1}(t) d t$,
$\left(y_{1 n}^{0}, v_{1 n}\right)_{\Omega} \varphi_{1}(0) \longrightarrow\left(y_{1}^{0}, v_{1}\right)_{\Omega} \varphi_{1}(0)$
and
$\int_{0}^{T}\left(y_{2 n}, v_{2 n}\right) \varphi_{2}^{\prime}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2 n}, v_{2 n}\right)+\right.$
$\left(b_{2}(t) y_{2 n}, v_{2 n}\right)_{\Omega}+$
$\left.\left(b(t) y_{1 n}, v_{2 n}\right)_{\Omega}\right] \varphi_{2}(t) d t \rightarrow$
$\int_{0}^{T}\left(y_{2}, v_{2}\right) \varphi_{2}^{\prime}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2}, v_{2}\right)+\right.$
$\left.\left(b_{2}(t) y_{2}, v_{2}\right)_{\Omega}+\left(b(t) y_{1}, v_{2}\right)_{\Omega}\right] \varphi_{2}(t) d t$
$\left(y_{2 n}^{0}, v_{2 n}\right)_{\Omega} \varphi_{2}(0) \rightarrow\left(y_{2}^{0}, v_{2}\right)_{\Omega} \varphi_{2}(0)$
On the other hand, let $w_{i n}=v_{i n} \varphi_{i}$ and $w_{i}=v_{i} \varphi_{i}$ then $\forall i=1,2, w_{i n} \rightarrow w_{i}$ strongly in $L^{2}(Q)$ and then $w_{\text {in }}$ is measurable w.r.t. ( $x, t$ ), using assumption (A-i), then applying Proposition (3.1), the integral $\int_{Q} f_{i}\left(x, t, y_{i n}\right) w_{i n} d x d t$ is continuous w.r.t. $\left(y_{i n}, w_{i n}\right)$, but $y_{i n} \rightarrow y_{i}$ strongly in $L^{2}(Q)$, then
$\int_{0}^{T}\left(f_{i}\left(y_{i n}\right), v_{i n}\right)_{\Omega} \varphi_{i}(t) d t \rightarrow$
$\int_{0}^{T}\left(f_{i}\left(y_{i}\right), v_{i}\right)_{\Omega} \varphi_{i}(t) d t, \forall i=1,2$

From (20-23) and the above converges, (18) and (19) become
$-\int_{0}^{T}\left(y_{1}, v_{1}\right) \varphi_{1}^{\prime}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1}, v_{1}\right)+\right.$ $\left.\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-\left(b(t) y_{2}, v_{1}\right)_{\Omega}\right] \varphi_{1}(t) d t=$ $\int_{0}^{T}\left(f_{1}\left(y_{1}\right), v_{1}\right)_{\Omega} \varphi_{1}(t) d t+$
$\int_{0}^{T}\left(u_{1}, v_{1}\right)_{\Gamma} \varphi_{1}(t) d t+\left(y_{1}^{0}, v_{1}\right)_{\Omega} \varphi_{1}(0)$
and
$-\int_{0}^{T}\left(y_{2}, v_{2}\right) \varphi_{2}^{\prime}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2}, v_{2}\right)+\right.$ $\left.\left(b_{2}(t) y_{2}, v_{2}\right)_{\Omega}+\left(b(t) y_{1}, v_{2}\right)_{\Omega}\right] \varphi_{2}(t) d t=$ $\int_{0}^{T}\left(f_{2}\left(y_{2}\right), v_{2}\right)_{\Omega} \varphi_{2}(t) d t+$
$\int_{0}^{T}\left(u_{2}, v_{2}\right)_{\Gamma} \varphi_{2}(t) d t+\left(y_{2}^{0}, v_{2}\right)_{\Omega} \varphi_{2}(0)$,
Now we have following two cases:

## Case 1:

Choose $\quad \varphi_{i} \in D[0, T], \quad$ i.e. $\quad \varphi_{i}(0)=$ $\varphi_{i}(T)=0, \forall i=1,2$ substituting these values for $\varphi_{i}$ in (24) and (25), finally using integration by parts for the $1^{\text {st }}$ terms in the L.H.S. of each one of the obtained equations, yield
$\int_{0}^{T}\left(y_{1 t}, v_{1}\right) \varphi_{1}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1}, v_{1}\right)+\right.$ $\left.\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-\left(b(t) y_{2}, v_{1}\right)_{\Omega} \varphi_{1}(t)\right] d t=$ $\int_{0}^{T}\left(f_{1}\left(y_{1}\right), v_{1}\right)_{\Omega} \varphi_{1}(t) d t+$
$\int_{0}^{T}\left(u_{1}, v_{1}\right)_{\Gamma} \varphi_{1}(t) d t$
and
$\int_{0}^{T}\left(y_{2 t}, v_{2}\right) \varphi_{2}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2}, v_{2}\right)+\right.$ $\left.\left(b_{2}(t) y_{2}, v_{2}\right)_{\Omega}+\left(b(t) y_{1}, v_{2}\right)_{\Omega} \varphi_{2}(t)\right] d t=$ $\int_{0}^{T}\left(f_{2}\left(y_{2}\right), v_{2}\right)_{\Omega} \varphi_{2}(t) d t+$
$\int_{0}^{T}\left(u_{2}, v_{2}\right)_{\Omega} \varphi_{2}(t) d t$,
i.e., $y_{1} \& y_{2}$ are solutions of the state equations (10a) \& (11a) respectively.

## Case 2:

Choose $\varphi_{i} \in C^{1}[0, T], \forall i=1,2$, such that $\varphi_{i}(T)=0 \& \varphi_{i}(0) \neq 0$,

Using integration by parts for the $1^{\text {st }}$ term in the L.H.S. of (26) \& (27)and subtracting two each obtained equations of (24),(25) respectively one gets
$\left(y_{1}^{0}, v_{1}\right)_{\Omega} \varphi_{1}(0)=\left(y_{1}(0), v_{1}\right)_{\Omega} \varphi_{1}(0), \Rightarrow$
$\left(y_{1}^{0}, v_{1}\right)_{\Omega}=\left(y_{1}(0), v_{1}\right)_{\Omega}$
i.e., the initial condition (10b) holds. Easily one can see that the initial condition (11b) holds.

The strong convergence for $\vec{y}_{n}$ in $L^{2}(I, V)$ :

By substituting $v_{1}=y_{1}$ and $v_{1}=y_{1 n}$ in (10a) and (13a) respectively and also substituting $v_{2}=y_{2}$ and $v_{2}=y_{2 n}$ in (11a) and (14a) respectively, integrating these four equations from $t=0$ to $t=T$ finally adding the equations which is obtained from (10a) with that obtained from (13a) to gather and the same thing happened for (11a), (14a), to get:
$\int_{0}^{T}\left\langle\vec{y}_{t}, \vec{y}\right\rangle d t+$
$\int_{0}^{T} c(t, \vec{y}, \vec{y}) d t=\int_{0}^{T}\left[\left(f_{1}\left(y_{1}\right), y_{1}\right)_{\Omega}+\right.$
$\left.\left(f_{2}\left(y_{2}\right), y_{2}\right)_{\Omega}\right] d t+\int_{0}^{T}\left(u_{1}, y_{1}\right)_{\Gamma}+\int_{0}^{T}\left(u_{2}, y_{2}\right)_{\Gamma}$
and
$\int_{0}^{T}\left\langle\vec{y}_{n t}, \vec{y}_{n}\right\rangle d t+\int_{0}^{T} c\left(t, \vec{y}_{n}, \vec{y}_{n}\right) d t=$
$\int_{0}^{T}\left[\left(f_{1}\left(y_{1 n}\right), y_{1 n}\right)+\left(f_{2}\left(y_{2 n}\right), y_{2 n}\right)\right] d t+$
$\int_{0}^{T}\left(u_{1}, y_{1 n}\right)_{\Gamma} d t+\int_{0}^{T}\left(u_{2}, y_{2 n}\right)_{\Gamma} d t$ $\qquad$
Using Lemma 1.2 in [13] for the $1^{\text {st }}$ terms in the L.H.S. of ( $28 \mathrm{a} \& \mathrm{~b}$ ), once get, $\frac{1}{2}\|\vec{y}(T)\|_{0}^{2}-\frac{1}{2}\|\vec{y}(0)\|_{0}^{2}+\int_{0}^{T} c(t, \vec{y}, \vec{y}) d t=$ $\int_{0}^{T}\left[\left(f_{1}\left(y_{1}\right), y_{1}\right)_{\Omega}+\left(f_{2}\left(y_{2}\right), y_{2}\right)_{\Omega}\right] d t+$ $\int_{0}^{T}\left(u_{1}, y_{1}\right)_{\Gamma} d t+\int_{0}^{T}\left(u_{2}, y_{2}\right)_{\Gamma} d t$
and
$\frac{1}{2}\left\|\vec{y}_{n}(T)\right\|_{0}^{2}-\frac{1}{2}\left\|\vec{y}_{n}(0)\right\|_{0}^{2}+\int_{0}^{T} c\left(t, \vec{y}_{n}, \vec{y}_{n}\right) d t=$ $\int_{0}^{T}\left[\left(f_{1}\left(y_{1 n}\right), y_{1 n}\right)_{\Omega}+\left(f_{2}\left(y_{2 n}\right), y_{2 n}\right)_{\Omega}\right] d t+$ $\int_{0}^{T}\left(u_{1}, y_{1 n}\right)_{\Gamma} d t+\int_{0}^{T}\left(u_{2}, y_{2 n}\right)_{\Gamma} d t$,
Now, consider the following equality:
$\frac{1}{2}\left\|\vec{y}_{n}(T)-\vec{y}(T)\right\|_{0}^{2}-\frac{1}{2}\left\|\vec{y}_{n}(0)-\vec{y}(0)\right\|_{0}^{2}+$
$\int_{0}^{T} c\left(t, \vec{y}_{n}-\vec{y}, \vec{y}_{n}-\vec{y}\right) d t=A_{1}-A_{2}-A_{3}$
where:
$\left(A_{1}\right)=\frac{1}{2}\left\|\vec{y}_{n}(T)\right\|_{0}^{2}-\frac{1}{2}\left\|\vec{y}_{n}(0)\right\|_{0}^{2}+$
$\int_{0}^{T} c\left(t, \vec{y}_{n}(T), \vec{y}_{n}(T)\right) d t$
$\left(A_{2}\right)=\frac{1}{2}\left(\vec{y}_{n}(T), \vec{y}(T)\right)-\frac{1}{2}\left(\vec{y}_{n}(0), \vec{y}(0)\right)+$
$\int_{0}^{T} c\left(t, \vec{y}_{n}(T), \vec{y}(T)\right) d t$
and
$\left(A_{3}\right)=\frac{1}{2}\left(\vec{y}(T), \vec{y}_{n}(T)-\vec{y}(T)\right)-\frac{1}{2}\left(\vec{y}(0), \vec{y}_{n}(0)-\right.$
$y(0))+\int_{0}^{T} c\left(t, \vec{y}(T), \vec{y}_{n}(T)-\vec{y}(T)\right) d t$
Since (form the projection theorem, see prove theorem 3.2)
$\vec{y}_{n}^{0}=\vec{y}_{n}(0) \rightarrow \vec{y}^{0}=\vec{y}(0)$ strongly in $\left(L^{2}(\Omega)\right)^{2}$
$\vec{y}_{n}(T) \rightarrow \vec{y}(T)$ strongly in $\left(L^{2}(\Omega)\right)^{2} \ldots .$. (31b)
Then
$\left(\vec{y}(0), \vec{y}_{n}(0)-\vec{y}(0)\right) \longrightarrow 0 \&\left(\vec{y}(T), \vec{y}_{n}(T)-\right.$ $\vec{y}(T)) \longrightarrow 0$
$\left\|\vec{y}_{n}(0)-\vec{y}(0)\right\|_{0}^{2} \rightarrow 0 \&\left\|\vec{y}_{n}(T)-\vec{y}(T)\right\|_{0}^{2} \rightarrow 0$
and since $\vec{y}_{n} \longrightarrow \vec{y}$ weakly in $\left(L^{2}(I, V)\right)^{2}$, then $\int_{0}^{T} c\left(t, \vec{y}(T), \vec{y}_{n}(T)-\vec{y}(T)\right) d t \rightarrow 0$

From proposition (3.1), the integral $\int_{0}^{T}\left(f_{i}\left(y_{i n}\right), y_{i n}\right) d t$ is continuous w.r.t. $y_{i}$, then
$\int_{0}^{T}\left[\left(f_{1}\left(y_{1 n}\right), y_{1 n}\right)+\left(f_{2}\left(y_{2 n}\right), y_{2 n}\right)\right] d t \rightarrow$
$\int_{0}^{T}\left[\left(f_{1}\left(y_{1}\right), y_{1}\right)+\left(f_{2}\left(y_{2}\right), y_{2}\right)\right] d t$
since $y_{\text {in }} \rightarrow y_{i}$ strongly in $L^{2}(Q), \forall i=1,2$.
Now, when $n \rightarrow \infty$ in both sides of (30), one has the following results:

1. The first two terms in the L.H.S. of (30) are tending to zero (from 31d)
2.Eq. $\left(A_{1}\right) \stackrel{\text { from }}{=} \int_{(29 \mathrm{~b})}^{T}\left[\left(f_{1}\left(y_{1 n}\right), y_{1 n}\right)+\right.$
$\left.\left(f_{2}\left(y_{2 n}\right), y_{2 n}\right)\right] d t+\int_{0}^{T}\left[\left(u_{1}, y_{1 n}\right)_{\Gamma}+\left(u_{2}, y_{2 n}\right)_{\Gamma}\right] d t$ $\underset{(31 \mathrm{f})}{\text { from }} \int_{0}^{T}\left[\left(f_{1}\left(y_{1}\right), y_{1}\right)_{\Omega}+\left(f_{2}\left(y_{2}\right), y_{2}\right)_{\Omega}\right] d t$ $+\int_{0}^{T}\left(u_{1}, y_{1}\right)_{\Gamma} d t+\int_{0}^{T}\left(u_{2}, y_{2}\right)_{\Gamma} d t$
2. Eq. $\left(A_{2}\right)$
$\rightarrow$ L.H.S.of $(29 \mathrm{a})=\int_{0}^{T}\left[\left(f_{1}\left(y_{1}\right), y_{1}\right)_{\Omega}+\right.$
$\left.\left(f_{2}\left(y_{2}\right), y_{2}\right)_{\Omega}\right] d t+\int_{0}^{T}\left[\left(u_{1}, y_{1 n}\right)_{\Gamma}+\right.$ $\left.\left(u_{2}, y_{2 n}\right)_{\Gamma}\right] d t$
3. The three terms in $\left(A_{3}\right)$ are tending to zero from(31c) and $\mathrm{c}(31 \mathrm{e})$.

From the above steps, (30) gives that:
$\int_{0}^{T} c\left(t, \vec{y}_{n}-\vec{y}, \vec{y}_{n}-\vec{y}\right) d t \rightarrow 0$
which means that:
$\bar{\alpha} \int_{0}^{T}\left\|\vec{y}_{n}-\vec{y}\right\|_{1}^{2} d t \rightarrow 0 \Rightarrow \vec{y}_{n} \rightarrow \vec{y}$ strongly in $\left(L^{2}(I, V)\right)^{2}$.

## Uniqueness of the solution:

Let $\vec{y}=\left(y_{1}, y_{2}\right), \overrightarrow{\hat{y}}=\left(\hat{y}_{1}, \hat{y}_{2}\right)$ be two solutions of the state equations (10a)-(11a), i.e. $\forall v_{1}, v_{2} \in V$, i.e., first from (10a), one has $\left\langle y_{1 t}, v_{1}\right\rangle+a_{1}\left(t, y_{1}, v_{1}\right)+\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-$ $\left(b(t) y_{2}, v_{1}\right)_{\Omega}=\left(f_{1}\left(y_{1}\right), v_{1}\right)_{\Omega}+\left(u_{1}, v_{1}\right)_{\Gamma}$, $\left\langle\hat{y}_{1 t}, v_{1}\right\rangle+a_{1}\left(t, \hat{y}_{1}, v_{1}\right)+\left(b_{1}(t) \hat{y}_{1}, v_{1}\right)_{\Omega}-$ $\left(b(t) \hat{y}_{2}, v_{1}\right)_{\Omega}=\left(f_{1}\left(\hat{y}_{1}\right), v_{1}\right)_{\Omega}+\left(u_{1}, v_{1}\right)_{\Gamma}$, By subtracting the $2^{\text {nd }}$ equation from the $1^{\text {st }}$ one, then substituting $v_{1}=y_{1}-\hat{y}_{1}$, to get
$\left\langle\left(y_{1}-\hat{y}_{1}\right)_{t}, y_{1}-\hat{y}_{1}\right\rangle+a_{1}\left(t, y_{1}-\hat{y}_{1}, y_{1}-\right.$
$\left.\hat{y}_{1}\right)+\left(b_{1}(t)\left(y_{1}-\hat{y}_{1}\right), y_{1}-\hat{y}_{1}\right)_{\Omega}-$
$\left(b(t)\left(y_{2}-\hat{y}_{2}\right), v_{1}\right)_{\Omega}=\left(f_{1}\left(y_{1}\right)-\right.$
$\left.f_{1}\left(\hat{y}_{1}\right), y_{1}-\hat{y}_{1}\right)_{\Omega}$,
Second form (11a), and by the same above way but for $y_{2}, \hat{y}_{2}$, the following equality is also obtained
$\left\langle\left(y_{2}-\hat{y}_{2}\right)_{t}, y_{2}-\hat{y}_{2}\right\rangle+a_{2}\left(t, y_{2}-\hat{y}_{2}, y_{2}-\right.$
$\left.\hat{y}_{2}\right)+\left(b_{2}(t)\left(y_{2}-\hat{y}_{2}\right), y_{2}-\hat{y}_{2}\right)_{\Omega}+$
$\left(b(t)\left(y_{1}-\hat{y}_{1}\right), y_{1}-\hat{y}_{1}\right)_{\Omega}=\left(f_{2}\left(y_{1}\right)-\right.$
$\left.f_{2}\left(\hat{y}_{1}\right), y_{2}-\hat{y}_{2}\right)_{\Omega}$,
Adding (32) and (33), applying Lemma 1.2 in [13] for the $1^{\text {st }}$ term of L.H.S of above equality, using assumption A-iii, yields
$\frac{1 d}{2 d t}\|\vec{y}-\overrightarrow{\hat{y}}\|_{0}^{2}+\bar{\alpha}\|\vec{y}-\overrightarrow{\hat{y}}\|_{1}^{2} \leq 1\left(f_{1}\left(y_{1}\right)-\right.$ $\left.f_{1}\left(\hat{y}_{1}\right), y_{1}-\hat{y}_{1}\right)_{\Omega}+\left(f_{2}\left(y_{2}\right)-f_{2}\left(\hat{y}_{2}\right), y_{2}-\right.$ $\left.\hat{y}_{2}\right)_{\Omega}$

Keep in mind the second term in the L.H.S. of (34) is positive, integration both sides of (34) with respect to $t$ from 0 to $t$, then using assumptions (A-ii) of the R.H.S, finally using Belman - Gronwall inequality, one gets $\|\vec{y}(t)-\overrightarrow{\hat{y}}(t)\|_{0}^{2}=0, \quad \forall t \Rightarrow\|\vec{y}-\overrightarrow{\hat{y}}\|_{L^{2}(I, V)}=$ $0 \Rightarrow \vec{y}=\overrightarrow{\hat{y}}$.

## 4. Existence of a classical Boundary Optimal Control

The following theorem and lemma are important to study the existence of a classical boundary optimal control vector.

## Theorem (4.1):

(a) In addition to assumptions (A), if $\vec{y}$ and $\vec{y}+\overrightarrow{\Delta y}$ are the states vectors corresponding to the controls vectors $\vec{u}$ and $\vec{u}+\overrightarrow{\Delta u}$, if $\vec{u}$ and $\overrightarrow{\Delta u}$ are bounded in $\left(L^{2}(\Sigma)\right)^{2}$, then
$\|\overrightarrow{\Delta y}\|_{L^{\infty}\left(1, L^{2}(\Omega)\right)} \leq K\|\overrightarrow{\Delta u}\|_{\Sigma^{\prime}}, \quad\|\overrightarrow{\Delta y}\|_{L^{2}(Q)} \leq$ $K\|\overrightarrow{\Delta u}\|_{\Sigma}$ and $\|\overrightarrow{\Delta y}\|_{L^{2}(I, V)} \leq K\|\overrightarrow{\Delta u}\|_{\Sigma}$
(b) With assumptions (A), the operator $\vec{u} \mapsto \vec{y}_{\vec{u}}$ from $\left(L^{2}(\Sigma)\right)^{2}$ in to $\left(L^{\infty}\left(I, L^{2}(\Omega)\right)\right)^{2}$, or in to $\left(L^{2}(I, V)\right)^{2}$, or in to $\left(L^{2}(Q)\right)^{2}$ is continuous.
Proof:
(a) Let $\vec{u}=\left(u_{1}, u_{2}\right) \in\left(L^{2}(\Sigma)\right)^{2}$ then by theorem (3.1) there exists $\vec{y}=\left(y_{1}=y_{u_{1}}, y_{2}=\right.$ $y_{u_{2}}$ ) which satisfies the weak forms (10a\&b) (11a\&b), $\forall v_{1}, v_{2} \in V$, also let $\overrightarrow{\hat{u}}=\left(\hat{u}_{1}, \hat{u}_{2}\right) \in$
$\left(L^{2}(\Sigma)\right)^{2}$, then $\overrightarrow{\hat{y}}=\left(\hat{y}_{1}, \hat{y}_{2}\right) \quad$ is the corresponding solution of the following weak forms
$\left\langle\hat{y}_{1 t}, v_{1}\right\rangle+a_{1}\left(t, \hat{y}_{1}, v_{1}\right)+\left(b_{1}(t) \hat{y}_{1}, v_{1}\right)_{\Omega}-$
$\left(b(t) \hat{y}_{2}, v_{1}\right)_{\Omega}=\left(f_{1}\left(\hat{y}_{1}\right), v_{1}\right)_{\Omega}+\left(\hat{u}_{1}, v_{1}\right)_{\Gamma}$
$\left(\hat{y}_{1}(0), v_{1}\right)_{\Omega}=\left(y_{1}^{0}, v_{1}\right)_{\Omega}$
and
$\left\langle\hat{y}_{2 t}, v_{2}\right\rangle+a_{2}\left(t, \hat{y}_{2}, v_{2}\right)+\left(b_{2}(t) \hat{y}_{2}, v_{2}\right)_{\Omega}+$ $\left(b(t) \hat{y}_{1}, v_{2}\right)_{\Omega}=\left(f_{1}\left(\hat{y}_{2}\right), v_{2}\right)_{\Omega}+\left(\hat{u}_{1}, v_{2}\right)_{\Gamma}$
$\left(\hat{y}_{2}(0), v_{2}\right)_{\Omega}=\left(y_{2}^{0}, v_{2}\right)_{\Omega}$
By subtracting ( $10 \mathrm{a} \& \mathrm{~b}$ ) and ( $11 \mathrm{a} \& \mathrm{~b}$ ) from (35a\&b), (36a\&b) respectively, setting $\Delta y_{1}=\hat{y}_{1}-y_{1}, \quad \Delta y_{2}=\hat{y}_{2}-y_{2}, \quad \Delta u_{1}=\hat{u}_{1}-$ $u_{1}$ and $\Delta u_{2}=\hat{u}_{2}-u_{2}$ in each one of the two obtained equations, i.e.
$\left\langle\Delta y_{1 t}, v_{1}\right\rangle+$
$a_{1}\left(t, \Delta y_{1}, v_{1}\right)+\left(b_{1}(t) \Delta y_{1}, v_{1}\right)_{\Omega}-$
$\left(b(t) \Delta y_{2}, v_{1}\right)_{\Omega}=\left(f_{1}\left(y_{1}+\Delta y_{1}\right), v_{1}\right)_{\Omega}-$
$\left(f_{1}\left(y_{1}\right), v_{1}\right)_{\Omega}+\left(\Delta u_{1}, v_{1}\right)_{\Gamma}$
$\left(\Delta y_{1}(0), v_{1}\right)_{\Omega}=0$
and
$\left\langle\Delta y_{2 t}, v_{2}\right\rangle+$
$a_{2}\left(t, \Delta y_{2}, v_{2}\right)+\left(b_{2}(t) \Delta y_{2}, v_{2}\right)_{\Omega}+$
$\left(b(t) \Delta y_{1}, v_{2}\right)_{\Omega}=\left(f_{2}\left(y_{2}+\Delta y_{2}\right), v_{2}\right)_{\Omega}$ $-\left(f_{2}\left(y_{2}\right), v_{2}\right)_{\Omega}+\left(\Delta u_{2}, v_{2}\right)_{\Gamma}$,
$\left(\Delta y_{2}(0), v_{2}\right)_{\Omega}=0$
By substituting $v_{1}=\Delta y_{1}$ in (37a) and $v_{2}=\Delta y_{2}$ in (38a), adding the obtained equations, using Lemma 1.2 in [13] for the $1^{\text {st }}$ term and Assumption (A-iii) in the L.H.S. of the obtained equation, one gets
$\frac{1 d}{2 d t}\|\overrightarrow{\Delta y}\|_{0}^{2}+\bar{\alpha}\|\overrightarrow{\Delta y}\|_{1}^{2} \leq \mid\left(f_{1}\left(y_{1}+\Delta y_{1}\right)-\right.$ $\left.f_{1}\left(y_{1}\right), \Delta y_{1}\right)\left|+\left|\left(f_{2}\left(y_{2}+\Delta y_{2}\right)-f_{2}\left(y_{2}\right), \Delta y_{2}\right)\right|\right.$ $\left|\left(\Delta u_{1}, \Delta y_{1}\right)\right|+\left|\left(\Delta u_{2}, \Delta y_{2}\right)\right|$
Since the $2^{\text {nd }}$ term of L.H.S. of (39) is positive, integrating both sides w.r.t. $t$ from 0 to $t$, then using assumptions (A-ii), and then using the Cauchy-Schwarz inequality for the R.H.S., finally using the Trace operator, to get $\|\overrightarrow{\Delta y}(t)\|_{0}^{2} \leq 4\|\overrightarrow{\Delta u}\|_{\Sigma}^{2}+L_{3} \int_{0}^{t}\|\overrightarrow{\Delta y}\|_{0}^{2} d t$,
where $L_{3}$ refers to a summation for constants
Applying the Belman-Gronwall inequality gives
$\|\overrightarrow{\Delta y}(t)\|_{0} \leq K\|\overrightarrow{\Delta u}\|_{Q}, t \in[0, T]$
$\Rightarrow\|\overrightarrow{\Delta y}\|_{L^{\infty}\left(I, L^{2}(\Omega)\right)} \leq K\|\overrightarrow{\Delta u}\|_{\Sigma^{\prime}}, t \in[0, T]$
From this result, easily once can get that
$\|\overrightarrow{\Delta y}\|_{L^{2}(Q)} \leq K\|\overrightarrow{\Delta u}\|_{\Sigma^{\prime}}$ and $\|\overrightarrow{\Delta y}\|_{L^{2}(I, V)} \leq K\|\overrightarrow{\Delta u}\|_{\Sigma}$
(b) Let $\overrightarrow{\Delta u}=\vec{u}_{1}-\vec{u}_{2}$ and $\overrightarrow{\Delta y}=\vec{y}_{1}-\vec{y}_{2}$ where $\vec{y}_{1}$ and $\vec{y}_{2}$ are the correspond states to the boundary controls $\vec{u}_{1}$ and $\vec{u}_{2}$, then frompart (a) of this theorem,once get that the operator $\vec{u} \mapsto \vec{y}$ is Lipschitz continuous from $\left(L^{2}(\Sigma)\right)^{2}$ in to $\left(L^{\infty}\left(I, L^{2}(\Omega)\right)\right)^{2}$. The other result is obtained easily.

## Assumptions (B):

Consider $g_{l i}$ and $h_{l i}$ (for each $l=0,1,2$ and $i=1,2$ ) is of Carathéodory type on $(Q \times \mathbb{R})$ and on $(\Sigma \times \mathbb{R})$ respectly and satisfies the following sub quadratic condition w.r.t. $y_{i}$ and $u_{i}$
$\left|g_{\ell i}\left(x, t, y_{i}\right)\right| \leq$
$\gamma_{l i}(x, t)+c_{l i}\left(y_{i}\right)^{2},\left|h_{l i}\left(x, t, u_{i}\right)\right| \leq \delta_{l i}(x, t)+$ $d_{l i}\left(u_{i}\right)^{2}$
where $y_{i}, u_{i} \in \mathbb{R}$ with $\gamma_{l i} \in L^{1}(Q), \delta_{l i} \in L^{1}(\Sigma)$

## Lemma (4.2):

With assumptions (B), the functional $G_{l}(\vec{u})$, is continuous on $\left(L^{2}(\Sigma)\right)^{2}$ for each $l=0,1,2$.

## Proof:

From assumptions (B), with Using proposition (3.1), the integrals $\int_{Q} g_{l i}\left(x, t, y_{i}\right) d x d t$ and $\int_{\Sigma} h_{l i}\left(x, t, u_{i}\right) d \sigma$ are continuous on $L^{2}(Q)$ and $L^{2}(\Sigma)$ respectively for $\quad$ each $i=1,2, \quad l=0,1,2 \Rightarrow G_{l}(\vec{u}) \quad$ is continuous on $\left(L^{2}(\Sigma)\right)^{2}, \forall l=0,1,2$.

## Theorem (4.3):

In addition to assumptions (A), and (B), If the set of controls is of the form $\vec{W}=\vec{W}_{\vec{U}}$ with $\vec{U}$ is convex and compact, $\vec{W}_{A} \neq \emptyset$, if for each $i=1,2, G_{1}(\vec{u})$ is independent of $u_{i}, G_{0}(\vec{u})$ $\operatorname{and} G_{2}(\vec{u})$ are convex w.r.t. $u_{i}$ for fixed $\left(x, t, y_{i}\right)$ for each $i=1,2$. Then there exists a boundary optimal control vector.

## Proof:

Since the set $W_{i}$ is convex, closed and bounded for each $i=1,2$, then $W_{1} \times W_{2}$ is convex, closed and bounded, which gives $W_{1} \times W_{2}$ is weakly compact. Since $\vec{W}_{A} \neq \emptyset$, then there exists $\overrightarrow{\vec{u}} \in \vec{W}_{A}$ and there exists a minimum sequence $\left\{\vec{u}_{k}\right\}$ with $\vec{u}_{k} \in \vec{W}_{A}, \forall k$ such that, $\lim _{k \rightarrow \infty} G_{0}\left(\vec{u}_{k}\right)=\inf _{\overrightarrow{\vec{u}} \in \vec{W}_{A}} G_{0}(\overrightarrow{\vec{u}})$.

But $\vec{W}$ is weakly compact, there exists a subsequence of $\left\{\vec{u}_{k}\right\}$ say again $\left\{\vec{u}_{k}\right\}$ which converges weakly to some point $\vec{u}$ in $\vec{W}$, i.e. $\vec{u}_{k} \rightarrow \vec{u}$ weakly in $\left(L^{2}(\Sigma)\right)^{2}$, and $\left\|\vec{u}_{k}\right\|_{\Sigma} \leq c$, $\forall k$.
From theorem (3.2) for each boundary control $\vec{u}_{k}$, the state equation has a unique solution $\vec{y}_{k}=\vec{y}_{\vec{u}_{k}}$ and the norms $\left\|\vec{y}_{k}\right\|_{L^{\infty}\left(I, L^{2}(\Omega)\right)}$, $\left\|\vec{y}_{k}\right\|_{L^{2}(Q)}$ and $\left\|\vec{y}_{k}\right\|_{L^{2}(I, V)}$ are bounded, then by Alaoglu's theorem there exists a subsequence of $\left\{\vec{y}_{k}\right\}$ say again $\left\{\vec{y}_{k}\right\}$ which converges weakly to some point $\vec{y}$ w.r.t the above norms, i.e. $\vec{y}_{k} \rightarrow \vec{y}$ weakly in $\left(L^{\infty}\left(I, L^{2}(\Omega)\right)\right)^{2}$, in $\left(L^{2}(Q)\right)^{2}$, and in $\left(L^{2}(I, V)\right)^{2}$.
Also, from theorem (3.2), the norm $\left\|\vec{y}_{k}\right\|_{L^{2}\left(I, V^{*}\right)}$ is bounded and since
$\left(L^{2}(I, V)\right)^{2} \subset\left(L^{2}(Q)\right)^{2} \cong\left(\left(L^{2}(Q)\right)^{*}\right)^{2} \subset$ $\left(L^{2}\left(I, V^{*}\right)\right)^{2}$
Then by using the First Compactness Theorem [13], there exists a subsequence of $\left\{\vec{y}_{k}\right\}$ say again $\left\{\vec{y}_{k}\right\}$ such that $\vec{y}_{k} \rightarrow \vec{y}$ strongly in $\left(L^{2}(Q)\right)^{2}$.
Since for each $k, \quad y_{1 k}$ and $y_{2 k}$ are corresponding solutions to the controls $u_{1 k}$ and $u_{2 k}$, i.e.,
$\left\langle y_{1 k t}, v_{1}\right\rangle+a_{1}\left(t, y_{1 k} v_{1}\right)+\left(b_{1}(t) y_{1 k}, v_{1}\right)_{\Omega}-$ $\left(b(t) y_{2 k}, v_{1}\right)_{\Omega}=\left(f_{1}\left(x, t, y_{1 k}\right), v_{1}\right)_{\Omega}$
$+\left(u_{1 k}, v_{1}\right)_{\Gamma}$
and
$\left\langle y_{2 k t}, v_{2}\right\rangle+$
$a_{2}\left(t, y_{2 k}, v_{2}\right)+\left(b_{2}(t) y_{2 k}, v_{2}\right)_{\Omega}+$
$\left(b(t) y_{1 k}, v_{2}\right)_{\Omega}=\left(f_{2}\left(x, t, y_{2 k}\right), v_{2}\right)_{\Omega}$
$+\left(u_{2 k}, v_{2}\right)_{\Gamma}$
Let $\varphi_{i} \in C^{1}[I]$, with $\varphi_{i}(T)=0, \forall i=$ 1,2. Multiplying both sides of (40) and (41) by $\varphi_{1}(t)$ and $\varphi_{2}(t)$ respectively, and then integrating both sides w.r.t. $t$ from 0 to $T$, finally using integration by parts for the $1^{\text {st }}$ terms in the L.H.S. of the two obtain equations, to get
$-\int_{0}^{T}\left(y_{1 k}, v_{1}\right) \dot{\varphi}_{1}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1 k} v_{1}\right)+\right.$
$\left.\left(b_{1}(t) y_{1 k}, v_{1}\right)_{\Omega}-\left(b(t) y_{2 k}, v_{1}\right)_{\Omega}\right] \varphi_{1}(t) d t=$
$\int_{0}^{T}\left(f_{1}\left(x, t, y_{1 k}\right), v_{1}\right)_{\Omega} \varphi_{1}(t) d t+$
$\int_{0}^{T}\left(u_{1 k}, v_{1}\right)_{\Gamma} \varphi_{1}(t) d t+\left(y_{1 k}(0), v_{1}\right)_{\Omega} \varphi_{1}(0)$
and
$-\int_{0}^{T}\left(y_{2 k}, v_{2}\right) \dot{\varphi}_{2}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2 k} v_{2}\right)+\right.$ $\left.\left(b_{2}(t) y_{2 k}, v_{2}\right)_{\Omega}+\left(b(t) y_{1 k}, v_{1}\right)_{\Omega}\right] \varphi_{2}(t) d t=$

$$
\begin{align*}
& \int_{0}^{T}\left(f_{2}\left(x, t, y_{2 k}\right), v_{2}\right)_{\Omega} \varphi_{2}(t) d t+ \\
& \int_{0}^{T}\left(u_{2 k}, v_{2}\right)_{\Gamma} \varphi_{2}(t) d t+\left(y_{2 k}(0), v_{2}\right)_{\Omega} \varphi_{2}(0) \tag{43}
\end{align*}
$$

Since $\vec{y}_{k} \rightarrow \vec{y}$ weakly in $\left(L^{2}(Q)\right)^{2}$ and in $\left(L^{2}(I, V)\right)^{2}$, then
$-\int_{0}^{T}\left(y_{1 k}, v_{1}\right) \dot{\varphi}_{1}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1 k}, v_{1}\right)+\right.$
$\left.\left(b_{1}(t) y_{1 k}, v_{1}\right)_{\Omega}-\left(\mathrm{b}(\mathrm{t}) y_{2 k}, v_{1}\right)_{\Omega}\right] \varphi_{1}(t) d t \rightarrow$
$-\int_{0}^{T}\left(y_{1}, v_{1}\right) \dot{\varphi}_{1}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1}, v_{1}\right)+\right.$ $\left.\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-\left(b(t) y_{2}, v_{1}\right)_{\Omega}\right] \varphi_{1}(t) d t$
and
$-\int_{0}^{T}\left(y_{2 k}, v_{2}\right) \dot{\varphi}_{2}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2 k}, v_{2}\right)+\right.$
$\left.\left(b_{2}(t) y_{2 k}, v_{2}\right)_{\Omega}+\left(\mathrm{b}(\mathrm{t}) y_{1 k}, v_{2}\right)_{\Omega}\right] \varphi_{1}(t) d t \rightarrow$ $-\int_{0}^{T}\left(y_{2}, v_{2}\right) \dot{\varphi}_{2}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2}, v_{2}\right)+\right.$
$\left.\left(b_{2}(t) y_{2}, v_{2}\right)_{\Omega}+\left(b(t) y_{1}, v_{2}\right)_{\Omega}\right] \varphi_{2}(t) d t$
Since $y_{i k}(0), y_{2 k}(0)$ are bounded in $L^{2}(\Omega)$ and from Projection theorem, one gets that: $\left(y_{i k}(0), v_{i}\right)_{\Omega} \varphi_{1}(0) \longrightarrow\left(y_{i}^{0}, v_{i}\right)_{\Omega} \varphi_{i}(0), i=1,2$

Let $\forall i=1,2, w_{i}=v_{i} \varphi_{1}(t)$, then $w_{i}(x, t)$ is fixed for fixed $(x, t) \in Q$, and then $w_{i} \in$ $L^{\infty}(I, V) \subset L^{2}(I, V) \subset L^{2}(Q)$. Let $v_{i} \in C[\bar{\Omega}]$, then $w_{i} \in C[\overline{\mathrm{Q}}]$ is measurable w.r.t. $(x, t)$ and let $\bar{f}_{i}\left(y_{1 k}\right)=f_{i}\left(y_{i k}\right) w_{l}$, then $\bar{f}_{i}: Q \times \mathbb{R} \rightarrow \mathbb{R}$ is continuous with respect to $y_{i k}$ for fixed $(x, t) \in Q$, then
$\| \bar{f}_{i}\left(x, t, y_{i k}(X) \| \leq \eta_{i}\left|w_{i}\right|+c_{1}\left|y_{i k}\right|\left|w_{i}\right|=\right.$
$\bar{\eta}_{i}^{2}+\bar{c}_{2}\left\|y_{i k}\right\|^{2}$, where $\bar{\eta}_{i}^{2}=\frac{1}{2}\left(\eta_{i}^{2}+\bar{c}_{1}\left|w_{i}\right|^{2}\right)$
By applying proposition (3.1), the integral $\int_{Q} f_{i}\left(y_{i k}\right) w_{i} d x d t$ is continuous w.r.t. $y_{1 k}$ but $y_{i k} \rightarrow y_{i}$ strongly in $L^{2}(Q)$, then
$\int_{Q} f_{i}\left(y_{i k}\right) w_{i} d x d t \rightarrow \int_{Q} f_{i}\left(y_{i}\right) w_{i} d x d t$,
$\forall w_{i} \in C[\bar{Q}]$
(44c)
on the other hand, since $u_{i k} \rightarrow u_{i}$ weakly in $L^{2}(\Sigma), \forall i=1,2$, then
$\int_{\Gamma}\left(u_{i k}, v_{i}\right) \varphi_{i}(t) d \gamma d t \rightarrow \int_{\Gamma}\left(u_{i}, v_{i}\right) \varphi_{i}(t) d \gamma d t$
(44d)
Finally, using (44a, b, c\& d) and (45b) in (4243), once get that
$-\int_{0}^{T}\left(y_{1}, v_{1}\right) \dot{\varphi}_{1}(t) d t+\int_{0}^{T}\left[a_{1}\left(t, y_{1}, v_{1}\right)+\right.$
$\left.\left(b_{1}(t) y_{1}, v_{1}\right)_{\Omega}-\left(b(t) y_{2}, v_{1}\right)_{\Omega}\right] \varphi_{1}(t) d t=$
$\int_{0}^{T}\left(f_{1}\left(x, t, y_{1}\right), v_{1}\right)_{\Omega} \varphi_{1}(t) d t+$
$\int_{0}^{T}\left(u_{1}, v_{1}\right)_{\Gamma} \varphi_{1}(t) d t+\left(y_{1}^{0}, v_{1}\right)_{\Omega} \varphi_{1}(0)$
$-\int_{0}^{T}\left(y_{2}, v_{2}\right) \dot{\varphi}_{2}(t) d t+\int_{0}^{T}\left[a_{2}\left(t, y_{2}, v_{2}\right)+\right.$
$\left.\left(b_{2}(t) y_{2}, v_{2}\right)_{\Omega}+\left(b(t) y_{1}, v_{2}\right)_{\Omega}\right] \varphi_{2}(t) d t=$
$\int_{0}^{T}\left(f_{2}\left(x, t, y_{2}\right), v_{2}\right)_{\Omega} \varphi_{2}(t) d t+$
$\left.\int_{0}^{T}\left(u_{2}, v_{2}\right)_{\Gamma} \varphi_{2}(t)\right) d t+\left(y_{2}^{0}, v_{2}\right)_{\Omega} \varphi_{2}(0)$
Equations (46) and (47) are also hold for each $v_{i} \in V, \forall i=1,2$ (Since $C[\bar{\Omega}]$ is dense in $V)$.

Now, using the same steps which are used in Case 1 and Case 2 in the proof of theorem 3.1, once get that $y_{1}$ and $y_{2}$ are solutions of the weak form of the state equations.

Now, since $g_{1 i}$ is independent of $u_{i}$, for each $i=1$,2. i.e.,
$G_{1}\left(\vec{u}_{k}\right)=\int_{Q} g_{11}\left(x, t, y_{1 k}\right) d x d t+$ $\int_{Q} g_{12}\left(x, t, y_{2 k}\right) d x d t$

From the continuity of $g_{1 i}\left(x, t, y_{i k}\right)$ w.r.t. $y_{k}$ and the proof of Lemma (4.2), the integral $\int_{Q} g_{1 i}\left(x, t, y_{i k}\right) d x d t$ is continuous w.r.t. $y_{i k}$ and since $\vec{y}_{k} \rightarrow \vec{y}$ strongly in $\left(L^{2}(Q)\right)^{2}$, hence from proposition 3.1, one gets that

$$
G_{1}(\vec{u})=\lim _{k \rightarrow \infty} G_{1}\left(\vec{u}_{k}\right)=0
$$

Again since for each $i=1,2$ and $l=0,2$, $g_{l i}\left(x, t, y_{i k}\right)$ is continuous w.r.t $y_{i k}$, then from the proof of Lemma (4.2), one has $\int_{Q} g_{l i}\left(x, t, y_{i k}\right) d x d t \rightarrow \int_{Q} g_{l i}\left(x, t, y_{i}\right) d x d t$

From the hypotheses on $h_{l i}, h_{l i}\left(x, t, u_{i}\right)$ is weakly lower semi continuous w.r.t. $u_{i}$, for each $i=1,2$ and $l=0,2$, then from (48), one has
$\int_{Q} g_{l i}\left(x, t, y_{i}\right) d x d t+\int_{\Sigma} h_{l i}\left(x, t, u_{i}\right) d \sigma \leq$
$\lim _{k \rightarrow \infty} \inf \int_{\Sigma} h_{l i}\left(x, t, u_{i k}\right) d \sigma+\int_{Q} g_{l i}\left(x, t, y_{i}\right) d x d t$
$=\lim _{k \rightarrow \infty} \inf \int_{\Sigma}\left(h_{l i}\left(x, t, u_{i k}\right) d \sigma+\right.$
$\lim _{k \rightarrow \infty} \int_{Q}\left(g_{l i}\left(x, t, y_{i}\right)-g_{l i}\left(x, t, y_{i k}\right)\right) d x d t+$
$\lim _{k \rightarrow \infty} \int_{Q} g_{l i}\left(x, t, y_{i k}\right) d x d t$
$=\lim _{k \rightarrow \infty} \inf \int_{\Sigma} h_{l i}\left(x, t, u_{i k}\right) d \sigma+$
$\lim _{k \rightarrow \infty} \inf \int_{Q} g_{l i}\left(x, t, y_{i k}\right) d x d t$
$\Rightarrow G_{l}(\vec{u}) \leq \lim _{k \rightarrow \infty} \inf G_{l}\left(\vec{u}_{k}\right)$, (for each $\left.l=0,2\right)$
But $G_{2}\left(\vec{u}_{k}\right) \leq 0, \forall k$, then $G_{2}(\vec{u}) \leq 0$, and one gets that $\vec{u} \in \vec{W}_{A}$ and
$G_{0}(\vec{u}) \leq \lim _{k \rightarrow \infty} \inf G_{0}\left(\vec{u}_{k}\right)=\lim _{k \rightarrow \infty} G_{0}\left(\vec{u}_{k}\right)=$
$\inf _{\overrightarrow{\bar{u}} \in \vec{W}_{A}} G_{0}\left(\overrightarrow{\vec{u}}_{k}\right)$
$\Rightarrow G_{0}(\vec{u})=\min _{\vec{u} \in \vec{W}_{A}} G_{0}\left(\overrightarrow{\vec{u}}_{k}\right) \Rightarrow \vec{u} \quad$ is $\quad$ a classical boundary optimal control.

## 5. The Necessary Conditions for Optimality

This section concerns with the derivation of the Fréchet derivative under some suitable
assumptions, it concerns also with the proof of theorem of necessary conditions for optimality so as the theorem of sufficient conditions under some additional assumptions. Hence, the following assumption is very useful.

## Assumptions (C):

If $f_{i y_{i}}, g_{l_{i} y_{i}}$ and $h_{l_{i} u_{i}}(l=0,1,2 \& i=1,2)$ are of Carathéodory type on $Q \times(\mathbb{R}), Q \times(\mathbb{R})$ and on $\Sigma \times(\mathbb{R})$ respectively, such that
$(\mathbb{R}), \quad Q \times(\mathbb{R})$ and on $\Sigma \times(\mathbb{R})$ respectively, such that
$\left|f_{i y_{i}}\left(x, t, y_{i}\right)\right| \leq L_{i}$
$\left|g_{l_{i} y_{i}}\left(x, t, y_{i}\right)\right| \leq \zeta_{l i}(x, t)+e_{l i}\left|y_{i}\right|$
and
$\left|h_{l_{i} u_{i}}\left(x, t, u_{i}\right)\right| \leq \eta_{l i}(x, t)+f_{l i}\left|u_{i}\right|$
where $(x, t) \in Q, y_{i}, u_{i} \in \mathbb{R}, \zeta_{l i}(x, t) \in L^{2}(Q)$, $\eta_{l i}(x, t) \in L^{2}(\Sigma)$, and $e_{l i}, f_{l i}>0$.

## Theorem (5.1):

Dropping index $l$, the Hamiltonian $H$ which is defined by
$H(x, t, \vec{y}, \vec{z}, \vec{u})=\sum_{i=1}^{2}\left(z_{i} f_{i}\left(x, t, y_{i}\right)+\right.$
$\left.g_{i}\left(x, t, y_{i}\right)+h_{i}\left(x, t, u_{i}\right)\right)$
and the adjoint state $z_{i}=z_{i u}\left(\right.$ where $\left.\mathrm{y}_{\mathrm{i}}=\mathrm{y}_{\mathrm{ui}}\right)$ equation satisfies
$-z_{1 t}-\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(a_{i j}(x, t) \frac{\partial z_{1}}{\partial x_{i}}\right)+b_{1}(x, t) z_{1}+$
$b(x, t) z_{2}=z_{1} f_{y_{1}}\left(x, t, y_{1}\right)+g_{y_{1}}\left(x, t, y_{1}\right)$, inQ
$-z_{2 t}-\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(b_{i j}(x, t) \frac{\partial z_{2}}{\partial x_{i}}\right)+b_{2}(x, t) z_{2}-$
$b(x, t) z_{1}=z_{2} f_{y_{2}}\left(x, t, y_{2}\right)+g_{y_{2}}\left(x, t, y_{2}\right)$, in $Q$
$z_{1}(x, T)=0$, on $\Omega$
$z_{2}(x, T)=0, \quad$ on $\Omega$
$\frac{\partial z_{1}}{\partial n}=0, \quad$ on $\Sigma$
$\frac{\partial z_{2}}{\partial n}=0, \quad$ on $\Sigma$
Then the Fréchet derivative of $G$ is given by,
$\dot{G}(\vec{u}) \overrightarrow{\Delta u}=\int_{\Sigma}\binom{z_{1}+h_{u_{1}}}{z_{2}+h_{u_{2}}} \cdot\binom{\Delta u_{1}}{\Delta u_{2}} d \sigma$
$=\int_{\Sigma} H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \cdot \overrightarrow{\Delta u} d \sigma$
where:
$H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u})=\sum_{i=1}^{2}\left(z_{i}+h_{i u_{i}}\left(x, t, u_{i}\right)\right)$

## Proof:

The weak forms of the adjoint equations are given by

$$
\begin{align*}
& -\left\langle z_{1 t}, v_{1}\right\rangle+a_{1}\left(t, z_{1}, v_{1}\right)+\left(b_{1}(t) z_{1}, v_{1}\right)_{\Omega}+ \\
& \left(b(t) z_{2}, v_{1}\right)_{\Omega}=\left(z_{1} f_{1 y_{1}}, v_{1}\right)_{\Omega}+\left(g_{y_{1}}, v_{1}\right)_{\Omega} \tag{49}
\end{align*}
$$

$$
\begin{align*}
& -\left\langle z_{2 t}, v_{2}\right\rangle+a_{2}\left(t, z_{2}, v_{2}\right)+\left(b_{2}(t) z_{2}, v_{2}\right)_{\Omega}- \\
& \left(b(t) z_{1}, v_{2}\right)_{\Omega}=\left(z_{2} f_{2 y_{2}}, v_{2}\right)_{\Omega}+\left(g_{y_{2}}, v_{2}\right)_{\Omega}
\end{align*}
$$

These weak forms have a unique solution and this can be proved by the same way which is used in the proof of theorem 3.2.

Now, substituting $v_{1}=z_{1}$ in (37) and $v_{2}=z_{2}$ in (38), integrating both sides with respect to $t$ from 0 to $T$, then adding two obtained equations to get
$\int_{0}^{T}\left\langle\overrightarrow{\Delta y}_{t}, \vec{z}\right\rangle d t+\int_{0}^{T}\left[a_{1}\left(t, \Delta y_{1}, z_{1}\right)+\right.$ $\left(b_{1}(t) \Delta y_{1}, z_{1}\right)_{\Omega}-\left(b(t) \Delta y_{2}, z_{1}\right)_{\Omega}+$ $a_{2}\left(t, \Delta y_{2}, v_{2}\right)+\left(b_{2}(t) \Delta y_{2}, v_{2}\right)_{\Omega}+$ $\left.\left(b(t) \Delta y_{1}, v_{2}\right)_{\Omega}\right] d t=\int_{0}^{T}\left(f_{1}\left(y_{1}+\right.\right.$ $\left.\left.\Delta y_{1}\right), z_{1}\right)_{\Omega} d t-\int_{0}^{T}\left(f_{1}\left(y_{1}\right), z_{1}\right)_{\Omega} d t+$ $\int_{0}^{T}\left(\Delta u_{1}, z_{1}\right)_{\Gamma} d t+\int_{0}^{T}\left(f_{2}\left(y_{2}+\Delta y_{2}\right), z_{2}\right)_{\Omega} d t-$ $\int_{0}^{T}\left(f_{2}\left(y_{2}\right), z_{2}\right)_{\Omega} d t+\int_{0}^{T}\left(\Delta u_{2}, z_{2}\right)_{\Gamma} d t$

From assumption (A-ii), and proposition (3.2), the Fréchet derivative of $f_{i}$ exists for each $y_{i} \in L^{2}$, which gives after using the result of Theorem (4.1)
$\int_{0}^{T}\left(f_{i}\left(x, t, y_{i}+\Delta y_{i}\right)-f_{i}\left(x, t, y_{i}\right), z_{i}\right)_{\Omega} d t=$
$\int_{0}^{T}\left(f_{i y_{i}} \Delta y_{i}, z_{i}\right) d t+\varepsilon_{i 1}\left(\Delta y_{i}\right)\left\|\Delta y_{i}\right\|_{Q}$
$\varepsilon_{1}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma^{\prime}}$
where, $\varepsilon_{1}(\overrightarrow{\Delta u}) \rightarrow 0$ as $\|\overrightarrow{\Delta u}\|_{\Sigma} \rightarrow 0$
By substituting (52) in the R.H.S. of (51), one has that
$\int_{0}^{T}\left\langle\overrightarrow{\Delta y}_{t}, \vec{z}\right\rangle d t+\int_{0}^{T}\left[a_{1}\left(t, \Delta y_{1}, z_{1}\right)+\right.$
$\left(b_{1}(t) \Delta y_{1}, z_{1}\right)_{\Omega}-\left(b(t) \Delta y_{2}, z_{1}\right)_{\Omega}+$ $a_{2}\left(t, \Delta y_{2}, v_{2}\right)+\left(b_{2}(t) \Delta y_{2}, v_{2}\right)_{\Omega}+$
$\left.\left(b(t) \Delta y_{1}, v_{2}\right)_{\Omega}\right] d t=\int_{0}^{T}\left(f_{1 y_{1}} \Delta y_{1}, z_{1}\right)_{\Omega} d t+$ $\int_{0}^{T}\left(f_{2 y_{2}} \Delta y_{2}, z_{2}\right)_{\Omega} d t+\int_{0}^{T}\left(\Delta u_{1}, z_{1}\right)_{\Gamma} d t+$
$\int_{0}^{T}\left(\Delta u_{2}, z_{2}\right)_{\Gamma} d t+\varepsilon_{1}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma}$
Now, substituting $v_{1}=\Delta y_{1}$ and $v_{2}=\Delta y_{2}$ in the adjoint equations (49) and (50) respectively, integrating both sides with respect to $t$ from 0 to $T$, using integrating by part for the $1^{\text {st }}$ term of each obtained equation, finally adding these two equations, to get:
$\int_{0}^{T}\left\langle\overrightarrow{\Delta y}_{t}, \vec{z}\right\rangle d t+\int_{0}^{T}\left[a\left(t, z_{1}, \Delta y_{1}\right)+\right.$ $\left(b_{1}(t) z_{1}, \Delta y_{1}\right)_{\Omega}+\left(b(t) z_{2}, \Delta y_{1}\right)_{\Omega}+$ $a_{2}\left(t, z_{2}, \Delta y_{2}\right)+\left(b_{2}(t) z_{2}, \Delta y_{2}\right)_{\Omega}-$ $\left.\left(b(t) z_{1}, \Delta y_{2}\right)_{\Omega}\right] d t=\int_{0}^{T}\left(z_{1} f_{1 y_{1}}, \Delta y_{1}\right)_{\Omega} d t+$ $\int_{0}^{T}\left(g_{1 y_{1}}, \Delta y_{1}\right)_{\Omega} d t+$
$\int_{0}^{T}\left(z_{2} f_{2 y_{2}}, \Delta y_{2}\right)_{\Omega} d t+\int_{0}^{T}\left(g_{2 y_{2}}, \Delta y_{2}\right)_{\Omega} d t$
By subtracting (54) from (53), one gets
$\int_{0}^{T}\left(g_{1 y_{1}}, \Delta y_{1}\right)_{\Omega} d t+\int_{0}^{T}\left(g_{2 y_{2}}, \Delta y_{2}\right)_{\Omega} d t=$
$\int_{0}^{T}\left(\Delta u_{1}, z_{1}\right)_{\Gamma} d t+\int_{0}^{T}\left(\Delta u_{2}, z_{2}\right)_{\Gamma} d t+$
$\varepsilon_{1}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma}$
Now, let $G_{A}(\vec{u})=\int_{Q} k_{1}\left(x, t, y_{1}, y_{2}\right) d x d t$,
$G_{B}(\vec{u})=\int_{\Sigma} k_{2}\left(x, t, u_{1}, u_{2}\right) d \sigma$
where:
$k_{1}\left(x, t, y_{1}, y_{2}\right)=g_{1}\left(x, t, y_{1}\right)+g_{2}\left(x, t, y_{2}\right)$,
$k_{2}\left(x, t, u_{1}, u_{2}\right)=h_{1}\left(x, t, u_{1}\right)+h_{2}\left(x, t, u_{2}\right)$,
From the Fréchet derivative and the result of theorem (4.1), one has
$G_{A}(\vec{u}+\overrightarrow{\Delta u})-G_{A}(\vec{u})=\int_{Q}\left(k_{1 y_{1}} \Delta y_{1}+\right.$
$\left.k_{1 y_{2}} \Delta y_{2}\right) d x d t+\varepsilon_{2}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma}$
$G_{B}(\vec{u}+\overrightarrow{\Delta u})-G_{B}(\vec{u})=\int_{\Sigma}\left(k_{2 u_{1}} \Delta u_{1}+\right.$
$\left.k_{2 u_{2}} \Delta u_{2}\right) d \sigma+\varepsilon_{3}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma}$
Adding (56) and (57), to get
$G(\vec{u}+\overrightarrow{\Delta u})-G(\vec{u})=\int_{Q}\left(g_{1 y_{1}} \Delta y_{1}+\right.$
$\left.g_{2 y_{2}} \Delta y_{2}\right) d x d t+\int_{\Sigma}\left(h_{1 u_{1}} \Delta u_{1}+\right.$
$\left.h_{2 u_{2}} \Delta u_{2}\right) d \sigma+\varepsilon_{4}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma}$
By substituting (55) in (58), gives
$G(\vec{u}+\overrightarrow{\Delta u})-G(\vec{u})=\int_{\Sigma}\left(\Delta u_{1}, z_{1}\right) d \sigma+$
$\int_{\Sigma}\left(\Delta u_{2}, z_{2}\right) d \sigma+\int_{\Sigma}\left(h_{1 u_{1}} \Delta u_{1}+\right.$
$\left.h_{2 u_{2}} \Delta u_{2}\right) d \sigma+\varepsilon_{5}(\overrightarrow{\Delta u})\|\overrightarrow{\Delta u}\|_{\Sigma}$
where $\varepsilon_{5}(\overrightarrow{\Delta u}) \rightarrow 0$ as $\|\overrightarrow{\Delta u}\|_{\Sigma} \rightarrow 0$
Using proposition (3.2), the Fréchet derivative of $G$ is
$(\dot{G}(\vec{u}), \overrightarrow{\Delta u})=\int_{\Sigma}\binom{z_{1}+h_{1 u_{1}}}{z_{2}+h_{2 u_{2}}}\binom{\Delta u_{1}}{\Delta u_{2}} d \sigma$.

## Theorem (5.2) Necessary Conditions for Optimality (Multipliers Theorem):

If $\vec{u} \in \vec{W}_{A}$ is an optimal control, i.e., there exists multipliers $\lambda_{l} \in R, l=0,1,2 \quad$ with $\lambda_{0} \geq 0, \lambda_{2} \geq 0, \sum_{l=0}^{2}\left|\lambda_{l}\right|=1$ such that $\sum_{l=0}^{2} \lambda_{l} \dot{G}_{l}(\vec{u})(\vec{u}-\vec{u}) \geq 0, \forall \vec{u} \in \vec{W}$
and
$\lambda_{2} G_{2}(\vec{u})=0$ (Transversality condition).
The above relation is equivalent (59) to the following (weak) point wise minimum principle.
$H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \vec{u}=\min _{\vec{u} \in U} H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \vec{u}$ a.e on $\Sigma$
$H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u})=\sum_{i=1}^{2}\left(z_{i}+h_{i u_{i}}\left(x, t, u_{i}\right)\right)$
with $z_{i}=\sum_{l=0}^{2} \lambda_{l} z_{l i}$ and $h_{i}=\sum_{l=0}^{2} \lambda_{l} h_{l i}$,
(for $i=1,2$ ).

## Proof:

With assumptions (A),(B) and (C), the functional $G_{l}(\vec{u})$ and $\dot{G}_{l}(\vec{u})$ (for $\left.l=0,1,2\right)$ are continuous and liner w.r.t. $(\vec{u}-\vec{u})$, then $G_{l}(\vec{u})$ is $\rho$-differentiable at each $\vec{u} \in \vec{W}, \forall \rho$, Then by using the Kuhn-Tucker-Lagrange multipliers theorems [14], there exists multipliers $\quad \lambda_{l} \in \mathbb{R}, l=0,1,2$, with $\lambda_{0} \geq 0$, $\lambda_{2} \geq 0, \sum_{l=0}^{2}\left|\lambda_{l}\right|=1$ such that (60) and (61) are hold, i.e.
$\left(\lambda_{0} \dot{G}_{0}(\vec{u})+\lambda_{1} \dot{G}_{1}(\vec{u})+\lambda_{2} \dot{G}_{2}(\vec{u})\right) \cdot(\vec{u}-\vec{u}) \geq$
$0, \forall \vec{u} \in \vec{W}$
Applying Theorem (5.1), setting $\overrightarrow{\Delta u}=\vec{u}-\vec{u}$ and substituting the Fréchet derivative of $G_{l}$, for $l=0,1,2$ in (58), one has that
$\sum_{l=0}^{2} \sum_{i=1}^{2} \int_{\Sigma}\left(\lambda_{l}\left(z_{l i}+h_{l i u_{i}}\right)\left(u_{i}-u_{i}\right) d \sigma \geq 0\right.$
Let $z_{i}=\sum_{l=0}^{2} \lambda_{l} z_{l i}, h_{i}=\sum_{l=0}^{2} \lambda_{l} h_{l i}$, for each $i=1,2$
$\Rightarrow \int_{\Sigma} H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \cdot \overrightarrow{\Delta u} d \sigma \geq 0$
To prove that (62) is equivalent to (61)
Let
$\vec{W}_{\vec{u}}=\left\{\vec{u} \in\left(L^{2}(\Sigma, \mathbb{R})\right)^{2} \mid \vec{u}(x, t) \in \vec{U}\right.$ a. e.in $\left.\Sigma\right\}$, with $\vec{U} \subset \mathbb{R}^{2}$, let $\{\vec{u}\}$ be a dense sequence in $\vec{W}_{\vec{U}}, \mu$ is Lebesgue measure on $\Sigma$ and let $S \subset \Sigma$ be a measurable set such that
$\vec{u}(x, t)=\left\{\begin{array}{l}\vec{u}_{k}(x, t), \text { if }(x, t) \in S \\ \vec{u}(x, t), \text { if }(x, t) \notin S\end{array}\right.$
Therefore (66) becomes
$\int_{S} H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \cdot\left(\vec{u}_{k}-\vec{u}\right) d S \geq 0, \forall S$
Using theorem (3.1), to get
$H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) .\left(\vec{u}_{k}-\vec{u}\right) \geq 0$, a.e. in $\Sigma$,
which means this inequality is satisfied on the boundary $\Sigma$ of the region $Q$ except in a subset $\Sigma_{k}$ such that $\mu\left(\Sigma_{k}\right)=0, \forall k$, where $\mu$ is a Lebesgue measure, i.e. the it satisfies on the boundary $\Sigma$ except in the union of $\mathrm{U}_{k} \Sigma_{k}$ with $\mu\left(\mathrm{U}_{k} \Sigma_{k}\right)=0$, but $\left\{\overrightarrow{\vec{u}}_{k}\right\}$ is a dense sequence in the control set $\vec{W}$, then there exists $\overrightarrow{\vec{u}} \in \vec{W}$ such that
$H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \vec{u}=\min _{\vec{w} \in \vec{U}} H_{\vec{u}}(x, t, \vec{y}, \vec{z}, \vec{u}) \vec{u}$, a.e. in $\Sigma, \forall \overrightarrow{\vec{u}} \in \vec{W}$. The proof of the converse is obtained directly.

## 6. Sufficient Conditions for Optimality <br> Theorem (6.1) (Sufficient Conditions for Optimality):

In addition to the assumptions (A), (B) and (C), suppose that $\vec{W}=\vec{W}_{\vec{U}}$ is convex, with $\vec{U}$ convex, $f_{i}$ and $g_{1 i}$ are affine w.r.t. $y_{i}$ for each $(x, t)$ and $i=1,2, g_{0 i}$ and $g_{2 i}$ are convex w.r.t. $y_{i}, h_{0 i}$ and $h_{2 i}$ are convex with respect to $u_{i}$ for each ( $x, t$ ) and, $\forall i=1,2$. Then the necessary conditions in Theorem (5.2) with $\lambda_{0}>0$, are sufficient.

## Proof:

Suppose $\vec{u}$ is satisfied the K.T.L. condition, the Transversality condition and $\vec{u} \in \vec{W}_{A}$, i.e.
$\int_{\Sigma}\binom{z_{1}+h_{1 u 1}}{z_{2}+h_{2 u 2}} \cdot \overrightarrow{\Delta u} d \sigma \geq 0, \quad \forall \vec{u} \in \vec{W} \quad$ and $\lambda_{2} G_{2}(\vec{u})=0$,
where $z_{i}=\sum_{l=0}^{2} \lambda_{l} z_{l i}$ and $h_{i}=\sum_{l=0}^{2} \lambda_{l} h_{l i}$ (for $i=1,2$ ).
Let $G(\vec{u})=\sum_{l=0}^{2} \lambda_{l} G_{l}(\vec{u})$, then
$\dot{G}(\vec{u}) \cdot \overrightarrow{\Delta u}=\sum_{l=0}^{2} \lambda_{l} \dot{G}_{l}(\vec{u}) \cdot \overrightarrow{\Delta u}=$
$\lambda_{0} \int_{\Sigma} \sum_{i=1}^{2}\left(z_{0 i}+h_{0 i u_{i}}\right) \Delta u_{i} d \sigma+$
$\lambda_{1} \int_{\Sigma} \sum_{i=1}^{2}\left(z_{1 i}+h_{1 i u_{i}}\right) \Delta u_{i} d \sigma$
$\lambda_{2} \int_{\Sigma} \sum_{i=1}^{2}\left(z_{2 i}+h_{2 i u_{i}}\right) \Delta u_{i} d \sigma \geq 0$
since the functions $f_{1} \& f_{2}$ in the R.H.S. of the state equation (1) and (2) are affine with respect to $y_{1}$, and $y_{2}$ for $\operatorname{each}(x, t) \in Q$ respectively, i.e.,
$f_{1}\left(x, t, y_{1}\right)=f_{11}(x, t) y_{1}+f_{12}(x, t)$
$f_{2}\left(x, t, y_{2}\right)=f_{21}(x, t) y_{2}+f_{22}(x, t)$
Let $\vec{u}=\left(u_{1}, u_{2}\right) \& \overrightarrow{\vec{u}}=\left(\bar{u}_{1}, \bar{u}_{2}\right)$ are two given controls and then (by Theorem (3.2)), $\vec{y}=$ $\left(y_{u_{1}}, y_{u_{2}}\right)=\left(y_{1}, y_{2}\right) \quad \& \quad \overrightarrow{\bar{y}}=\left(\bar{y}_{\bar{u}_{1}}, \bar{y}_{\bar{u}_{2}}\right)=$ ( $\bar{y}_{1}, \bar{y}_{2}$ ) are their corresponding solutions, i.e. and for the $1^{\text {st }}$ state equations and their corresponding initial condition
$y_{1 t}-\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(a_{i j}(x, t) \frac{\partial y_{1}}{\partial x_{j}}\right)+b_{1}(x, t) y_{1}-$
$b(x, t) y_{2}=f_{11}(x, t) y_{1}+f_{12}(x, t)$
$\sum_{i, j=1}^{n} a_{i j} \frac{\partial y_{1}}{\partial n}=u_{1}(x, t)$,
$y_{1}(x, 0)=y_{1}^{0}(x)$
and
$\bar{y}_{1 t}-\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(a_{i j}(x, t) \frac{\partial \bar{y}_{1}}{\partial x_{j}}\right)+b_{2}(x, t) \bar{y}_{1}-$ $b(x, t) \bar{y}_{2}=f_{11}(x, t) \bar{y}_{1}+f_{12}(x, t)$
$\sum_{i, j=1}^{n} b_{i j} \frac{\partial y_{1}}{\partial n}=\bar{u}_{1}(x, t)$
$\bar{y}_{1}(x, 0)=y_{1}^{0}(x)$

By multiplying the $1^{\text {st }}$ above equation and its initial condition by $\theta, \theta \in[0,1]$, and the $2^{\text {nd }}$ equation and its initial condition by ( $1-\theta$ ), and adding the obtained equations and their obtained initial conditions, one has
$\left(\theta y_{1}+(1-\theta) \bar{y}_{1}\right)_{t}-$
$\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(a_{i j}(x, t) \frac{\partial\left(\theta y_{1}+(1-\theta) \bar{y}_{1}\right)}{\partial x_{j}}\right)+$
$b_{1}(x, t)\left(\theta y_{1}+(1-\theta) \bar{y}_{1}\right)-b(x, t)\left(\theta y_{2}+\right.$
$\left.(1-\theta) \bar{y}_{2}\right)=f_{11}(x, t)\left(\theta y_{1}+(1-\theta) \bar{y}_{1}\right)+$
$f_{12}(x, t)$
$\theta y_{1}(x, 0)+(1-\theta) \bar{y}_{1}(x, 0)=y_{1}^{0}(x) \ldots$. (63b)
$\sum_{i, j=1}^{n} a_{i j} \frac{\partial\left(\theta y_{1}+(1-\theta) \bar{y}_{1}\right)}{\partial n}=\left(\theta u_{1}+\left(1-\theta \bar{u}_{1}\right)\right.$ on $\Sigma$

By the same way and for the second differential equations and their initial conditions, one gets
$\left(\theta y_{2}+(1-\theta) \bar{y}_{2}\right)_{t}-$
$\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{i}}\left(b_{i j}(x, t) \frac{\partial\left(\theta y_{2}+(1-\theta) \bar{y}_{2}\right)}{\partial x_{j}}\right)+$
$b_{2}(x, t)\left(\theta y_{21}+(1-\theta) \bar{y}_{2}\right)+b(x, t)\left(\theta y_{1}+\right.$
$\left.(1-\theta) \bar{y}_{1}\right)=f_{21}(t)\left(\theta y_{2}+(1-\theta) \bar{y}_{2}\right)+$
$f_{22}(x, t)$
$\theta y_{2}(x, 0)+(1-\theta) \bar{y}_{2}(x, 0)=y_{2}^{0}(x) \ldots(64 b)$
$\sum_{i, j=1}^{n} b_{i j} \frac{\partial\left(\theta y_{2}+(1-\theta) \bar{y}_{2}\right)}{\partial n}=\left(\theta u_{2}+(1-\right.$
$\theta \bar{u}_{2}$ ),on $\Sigma$
Equations (63) and (64), tell us that the control $\overrightarrow{\tilde{u}}=\left(\tilde{u}_{1}, \tilde{u}_{2}\right)$, with $\overrightarrow{\tilde{u}}=\theta \vec{u}+(1-\theta) \overrightarrow{\vec{u}}$ has the corresponding solutions, $\overrightarrow{\tilde{y}}=\left(\tilde{y}_{1}, \tilde{y}_{2}\right)$, with $\overrightarrow{\tilde{y}}=\theta \vec{y}+(1-\theta) \overrightarrow{\vec{y}}$, which means the operator $\vec{u} \mapsto \vec{y}_{\vec{u}}$ is convex - linear with respect to $(\vec{y}, \vec{u})$ for $\operatorname{each}(x, t)$.
Now, since for each $i=1,2, \quad g_{1 i}\left(x, t, y_{i}\right)$ is affine w.r.t $y_{i}, \forall(x, t) \in Q$ and $h_{1 i}\left(x, t, u_{i}\right)$ is affine w.r.t. $u_{i}, \forall(x, t) \in \Sigma$, i.e.,
$g_{1 i}\left(x, t, y_{i}\right)=I_{1 i}(x, t) y_{i}+I_{2 i}(x, t)$,
$h_{1 i}\left(x, t, y_{i}\right)=I_{1 i}(x, t) u_{i}+I_{3 i}(x, t)$.
Let $\vec{u} \& \overrightarrow{\vec{u}}$ are two controls, and $\vec{y}=\vec{y}_{\vec{u}}$ \& $\overrightarrow{\bar{y}}=\overrightarrow{\bar{y}}_{\overrightarrow{\vec{u}}}$ are their corresponding solutions, then
$G_{1}(\theta \vec{u}+(1-\theta) \overrightarrow{\bar{u}})=$
$\sum_{i=1}^{2} \int_{Q} g_{1 i}\left(x, t, y_{i\left(\theta u_{i}+(1-\theta) \bar{u}_{i}\right)}\right) d x d t+$
$\sum_{i=1}^{2} \int_{\Sigma}\left[h_{1 i}\left(x, t, \theta u_{i}+(1-\theta) \bar{u}_{i}\right)\right] d \sigma$
$=\sum_{i=1}^{2}\left[\int_{Q} I_{1 i}\left(x, t, y_{i\left(\theta u_{i}+(1-\theta) \bar{u}_{i}\right)}\right)+\right.$
$\left.I_{2 i}(x, t)\right] d x d t+\sum_{i=1}^{2}\left[\int_{\Sigma} I_{1 i}\left(x, t, \theta u_{i}+\right.\right.$
$\left.\left.(1-\theta) \bar{u}_{i}\right)+I_{3 i}(x, t)\right] d \sigma$
Since the operator $\vec{u} \mapsto \vec{y}_{\vec{u}}$ is convex - linear, then
$G_{1}(\theta \vec{u}+(1-\theta) \overrightarrow{\vec{u}})=$
$\sum_{i=1}^{2}\left\{\int_{Q}\left[I_{1 i}(x, t)\left(\theta y_{i}+(1-\theta) \bar{y}_{i}\right)+\right.\right.$
$\left.I_{2 i}(x, t)\right] d x d t+$
$\sum_{i=1}^{2}\left\{\int_{\Sigma}\left[I_{1 i}(x, t)\left(\theta u_{i}+(1-\theta) \bar{u}_{i}\right)+\right.\right.$
$\left.I_{3 i}(x, t)\right] d \sigma$
$G_{1}(\theta \vec{u}+(1-\theta) \overrightarrow{\vec{u}})=\theta G_{1}(\vec{u})+$
$(1-\theta) G_{1}(\overrightarrow{\vec{u}})$
$\Rightarrow G_{1}(\vec{u})$ is convex - linear w.r.t. $(\vec{y}, \vec{u})$,
$\forall(x, t) \in Q$.
From the Assumptions on $g_{0 i} \& g_{2 i}\left(h_{0 i} \& h_{2 i}\right)$ $\forall i=1,2$, the integrals
$\sum_{i=1}^{2} \int_{Q} g_{0 i} d x d t \quad \& \quad \sum_{i=1}^{2} \int_{Q} g_{2 i} d x d t$ ( $\sum_{i=1}^{2} \int_{\Sigma} h_{0 i} d \sigma \& \sum_{i=1}^{2} \int_{\Sigma} h_{2 i} d \sigma$ ) are convex w.r.t. $y_{i} \forall(x, t) \in \mathrm{Q}$ and (are convex w.r.t. $u_{i}$ $\forall(x, t) \in \Sigma)$, then $G_{0}(\vec{u}) \& G_{2}(\vec{u})$ are convex w.r.t. $(\vec{y}, \vec{u}), \quad(\forall(x, t) \in Q, \forall(x, t) \in \Sigma)$, i.e. $G(\vec{u})$ is convex w.r.t. $(\vec{y}, \vec{u}), \quad(\forall(x, t) \in$ $Q, \forall(x, t) \in \Sigma)$. On the other hand, since $\vec{W}=\vec{W}_{\vec{U}}$ is convex, and the Fréchet derivative of $G_{l}(\vec{u}),(l=0,1,2)$ exists for each $\vec{u} \in \vec{W}$ and it is continuous (by Theorem (5.1) and assumptions (A),(B)and (C)), then it satisfies $\dot{G}(\vec{u}) \overrightarrow{\Delta u} \geq 0$, which means $G(\vec{u})$ has a minimum at $\vec{u}$, i.e.
$\lambda_{0} G_{0}(\vec{u})+\lambda_{1} G_{1}(\vec{u})+\lambda_{2} G_{2}(\vec{u}) \leq \lambda_{0} G_{0}(\vec{w})+$ $\lambda_{1} G_{1}(\vec{w})+\lambda_{2} G_{2}(\vec{w})$
Let $\vec{w} \in \vec{W}_{A}$, with $\lambda_{2} \geq 0$, then Transversality condition (64), gives
$\Rightarrow G_{0}(\vec{u}) \leq G_{0}(\vec{w}), \forall \vec{w} \in \vec{W}$, since $\left(\lambda_{0}>0\right)$
Hence $\vec{u}$ is a continuous classical boundary optimal control for the problem.

## Conclusions

In this paper, the existence and uniqueness theorem of a continuous classical boundary optimal control vector governing by the considered couple of nonlinear partial differential equation of parabolic type with equality and inequality constraints is proved using the Galerkin method, the existence of a classical boundary optimal control is proved under a suitable conditions, while the existence and uniqueness solution of the couple of adjoint vector equations associated with the considered couple equations of the state equations is proved and the derivation of the Fréchet derivative of the Hamiltonian is derived. Finally the theorem of necessary conditions and the theorem of sufficient
conditions of optimality problem with equality and inequality constrained are proved.

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